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# Casimir Effects, Riemann, Einstein, and Cosmology

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# Outline of this presentation

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# Einstein's Cosmological Constant

Our universe seems to be spatially flat and to possess a non-vanishing cosmological constant

- For cosmologists and general relativists: a great mistake why was it put there, in the first place? (Einstein)
- For elementary particle physicists: a great embarrassment no way to get rid off (Coleman, Weinberg, Polchinski)
- The cc  $\Lambda$  is indeed a peculiar quantity
  - has to do with cosmology Einstein's eqs., FRW universe
  - has to do with the local structure of elementary particle physics stress-energy density  $\mu$  of the vacuum

$$L_{cc} = \int d^4x \sqrt{-g} \mu^4 = \frac{1}{8\pi G} \int d^4x \sqrt{-g} \lambda$$

In other words: two contributions on the same footing (Zel'dovich, 68)

$$\frac{\Lambda c^2}{8\pi G} + \frac{1}{\text{Vol}} \frac{\hbar c}{2} \sum_i \omega_i$$



# Einstein Eqs, FLRW Sol, Hubble Const

Einstein Equations (1915-17):  $G_{\mu\nu} - \lambda g_{\mu\nu} = -(8\pi G/c^4)T_{\mu\nu}$

Geometry = Energy-Matter

$G_{\mu\nu}$  linear combination of the metric  $g_{\mu\nu}$  and 1st & 2nd derivatives

$T_{\mu\nu}$  energy-momentum tensor

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu} + \lambda g_{\mu\nu}$$

Schwarzschild solution (1916)

$r, \theta, \varphi$  comoving co

$$ds^2 = \left(1 - \frac{2MG}{r}\right) dt^2 - \left(1 - \frac{2MG}{r}\right)^{-1} dr^2 - r^2 d\theta^2 - r^2 \sin^2 \theta d\varphi^2$$

Friedmann-Lemaître-Robertson-Walker (1935-36) sol (A. Friedmann 1922)

$$ds^2 = dt^2 - a^2(t) \left( \frac{dr^2}{1 - kr^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\varphi^2 \right)$$

most general family: *homogeneous* and *isotropic*,  $k$  parameter  $\pm 1, 0$

One field eq looks like Newtonian eq for the gravit pot:  $\nabla^2 \phi = 4\pi G (\rho + 3p/c^2)$

density & pressure contribute to the gravit pot  $\lambda = 8\pi G \rho_{vac}$ ,  $p_{vac} = -\rho_{vac} c^2$

From the FRW metric and Einstein Eqs, an “equation of motion” of the universe

$$H^2 \equiv \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho + \frac{\lambda}{3} - \frac{k}{a^2}$$

# From GR to Cosmology

With the definitions:

$$\Omega_m \equiv \frac{8\pi G}{3H^2} \rho_m, \quad \Omega_\lambda \equiv \frac{\lambda}{3H^2}, \quad \Omega_k \equiv -\frac{k}{H^2}$$

The equation of motion becomes

$$\left(\frac{da}{dt}\right)^2 = H_0^2 \left[ \frac{\Omega_m^{(0)}}{a} + a^2 \Omega_\lambda^{(0)} + \Omega_k^{(0)} \right]$$

(the superscript (o) represents quantities measured at the present time)

In other terms, **Friedmann equation in Cosmology:**

$$\frac{\dot{a}^2}{a^2} = H_0^2 \left[ \Omega_R \left(\frac{a_0}{a}\right)^4 + \Omega_{NR} \left(\frac{a_0}{a}\right)^3 + \Omega_k \left(\frac{a_0}{a}\right)^2 + \Omega_\lambda \right]$$

$\Omega_R$  relativistic matter ( $p_R = \frac{1}{3}\rho_R$ ;  $\rho_R \propto a^{-4}$ )

Mach's princ

$\Omega_{NR}$  nonrelativistic matter ( $p_{NR} = 0$ ;  $\rho_{NR} \propto a^{-3}$ )

$\Omega_\lambda$  cosmological constant ( $p_\lambda = -\rho_\lambda$ ;  $\rho_\lambda = \text{const}$ )

$\Omega = \Omega_R + \Omega_{NR} + \Omega_\lambda$  total energy density (cosmic triangle)

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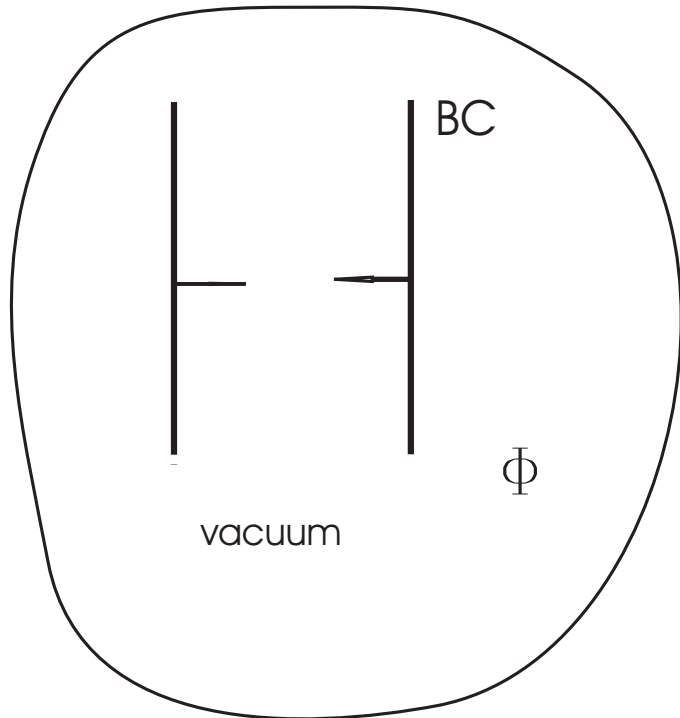
Even then: Has the final value real sense ?



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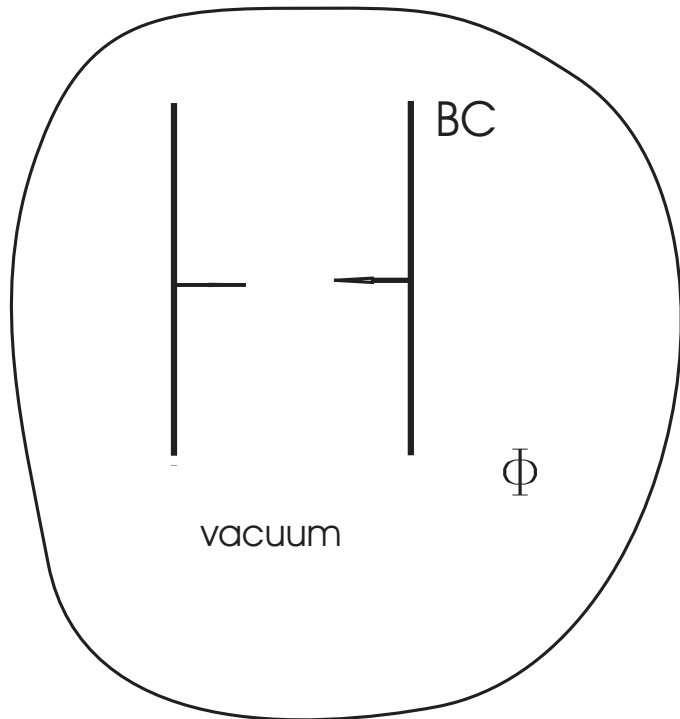
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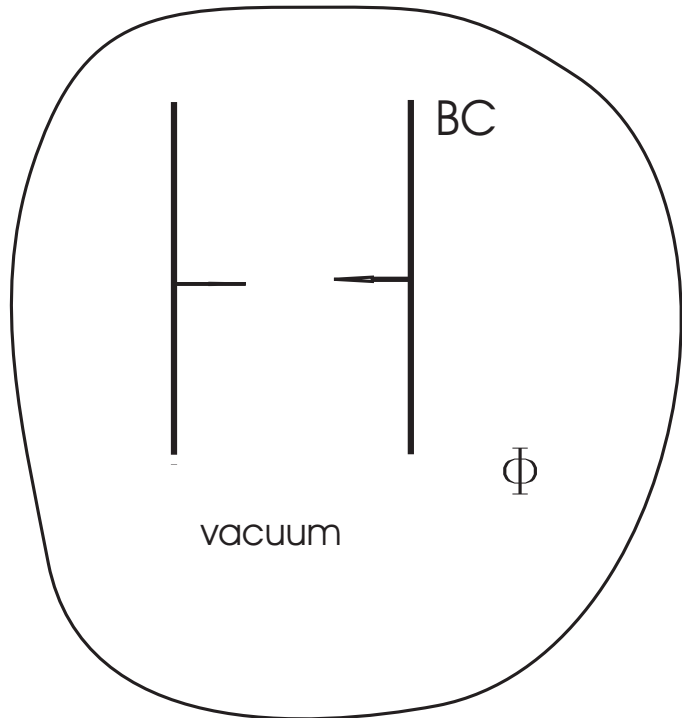
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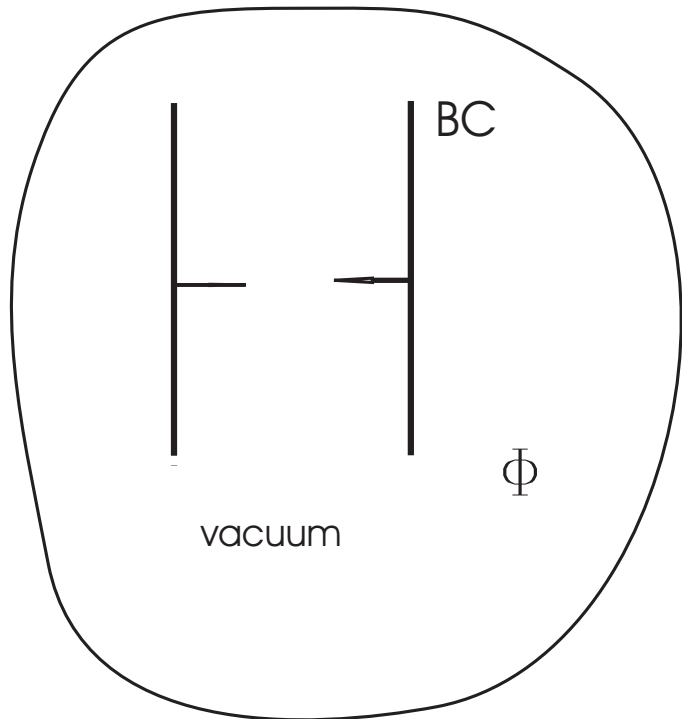
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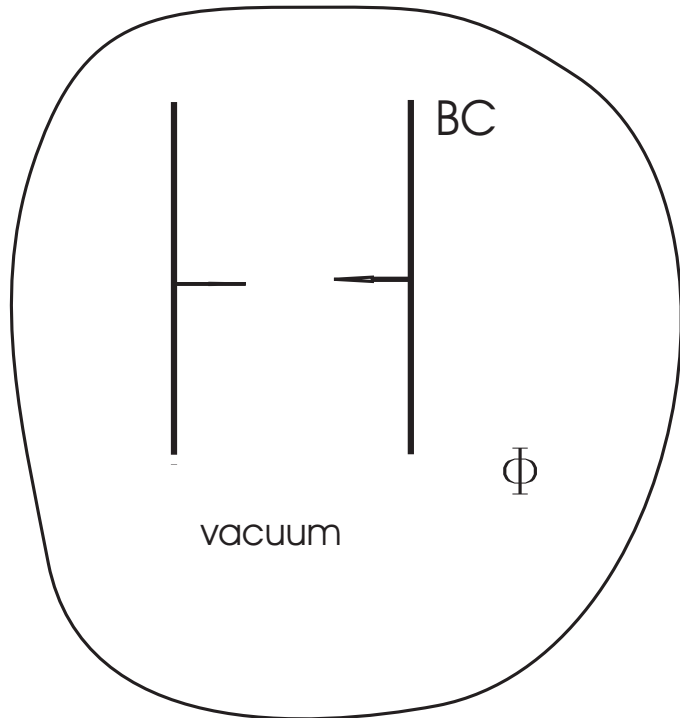
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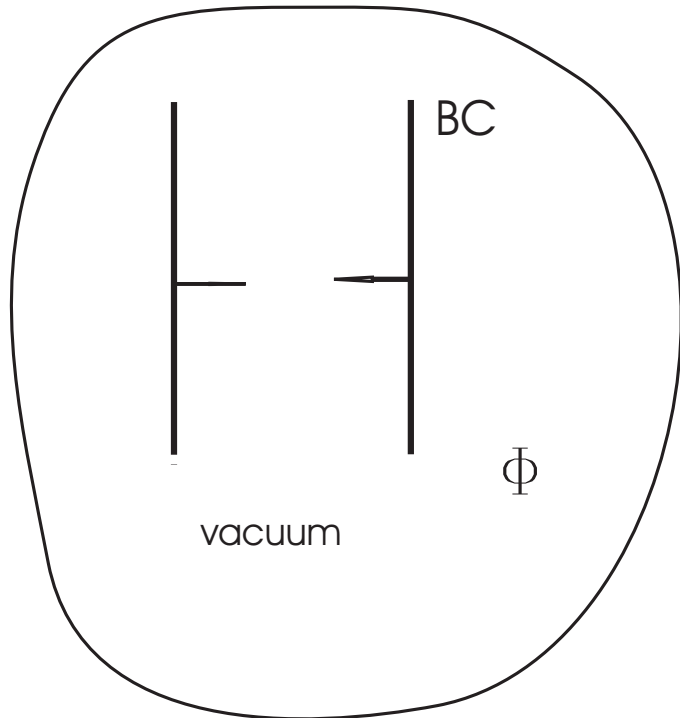
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Universal process:

- Sonoluminescence (Schwinger)
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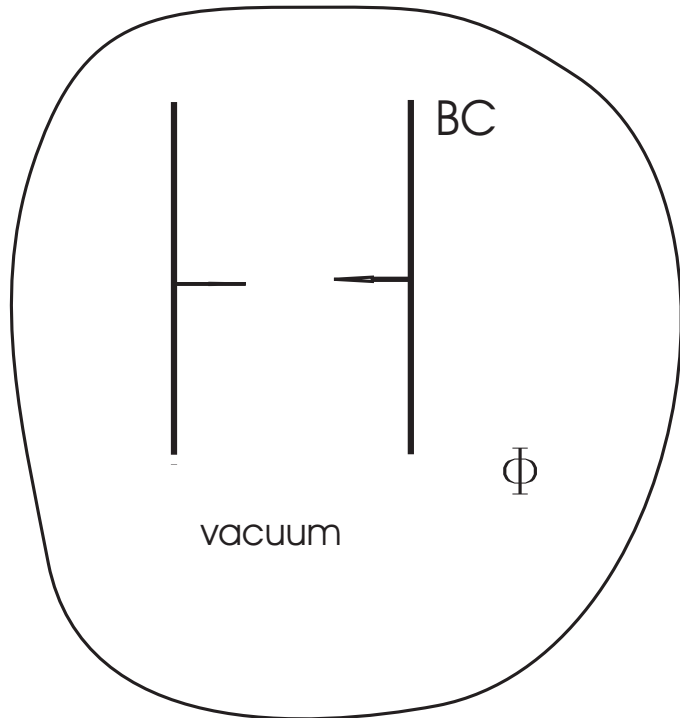
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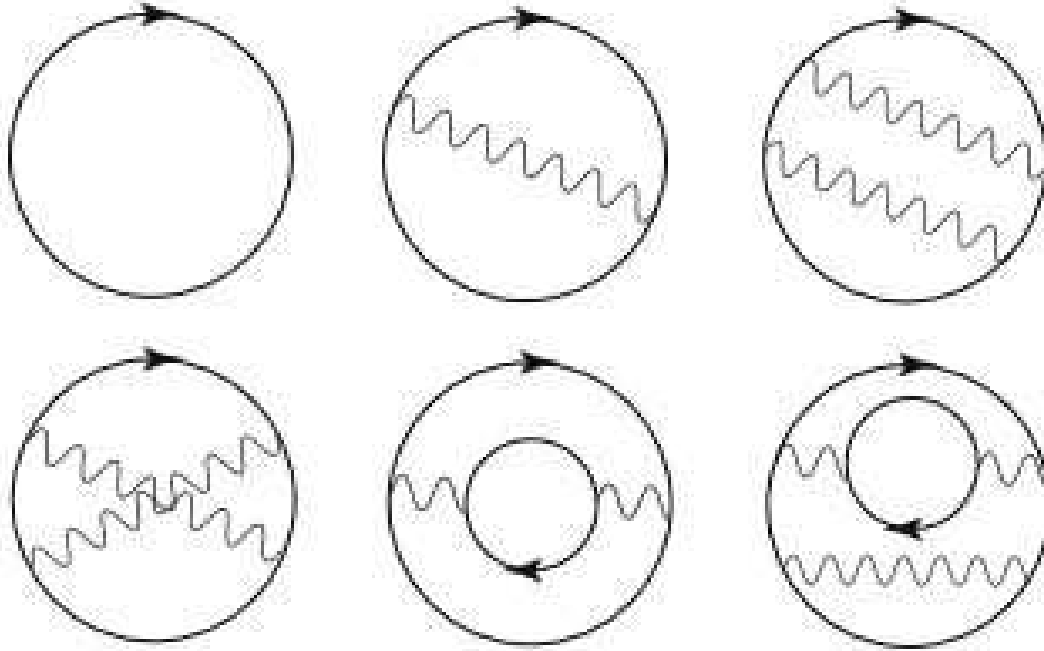
Van der Waals, Lifschitz theory

- Dynamical CE  $\Leftarrow$
- Lateral CE
- Extract energy from vacuum
- CE and the cosmological constant  $\Leftarrow$



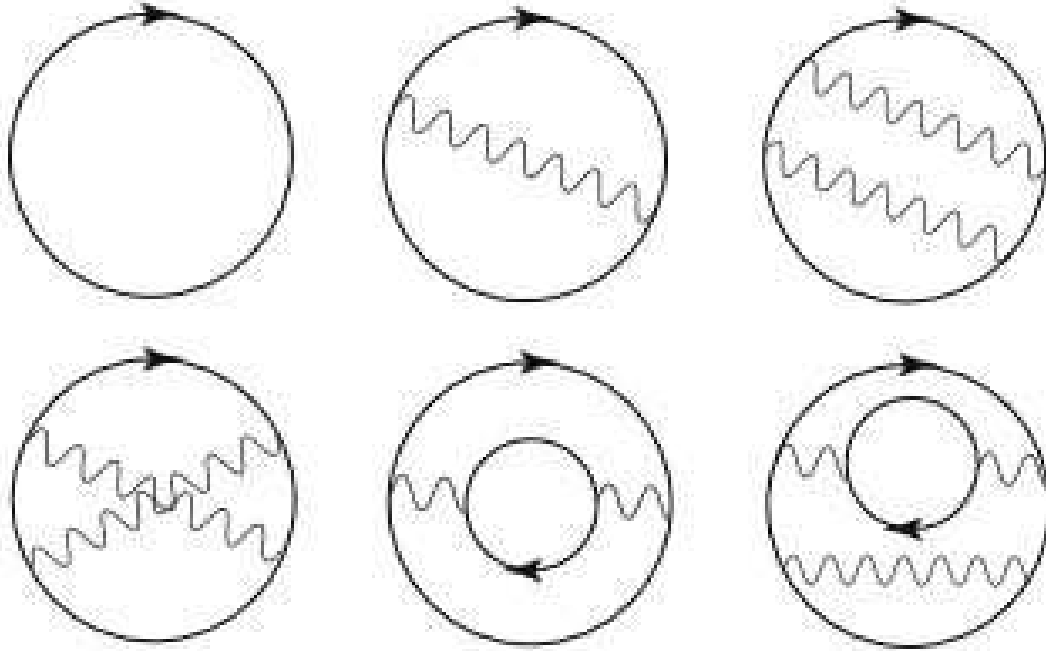
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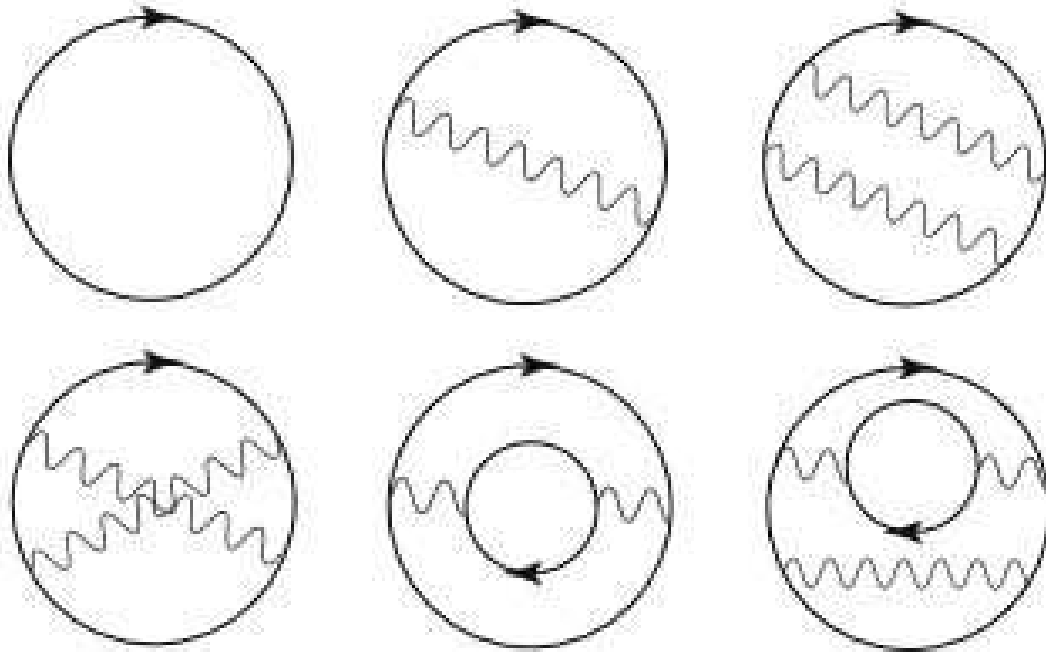
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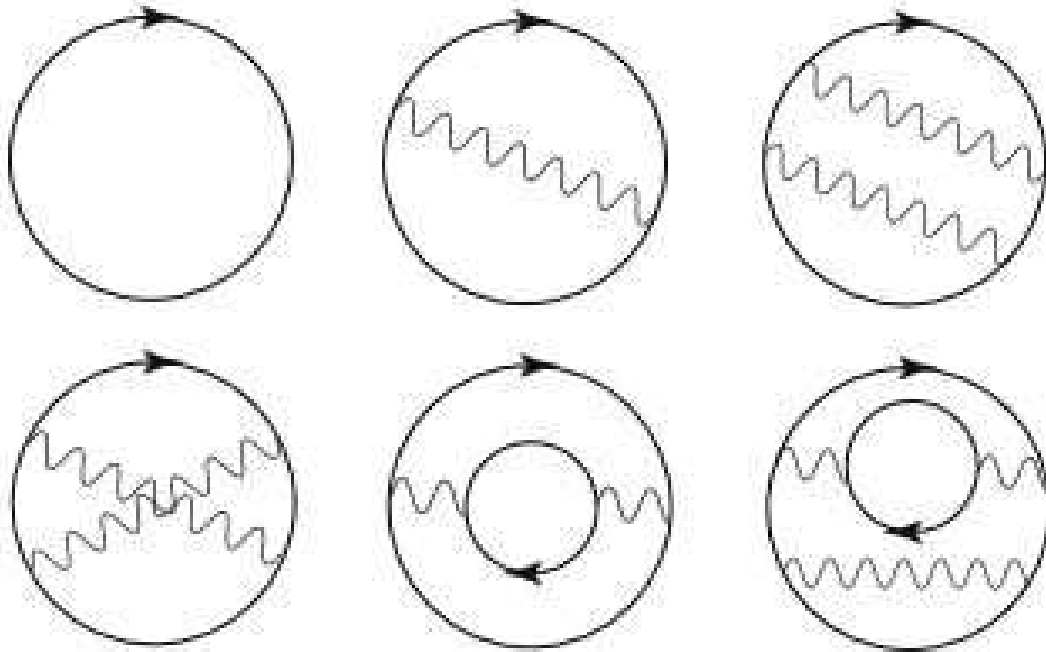
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$$\mathcal{E} = \frac{\hbar}{2\pi} \text{Im} \int d\omega \omega \text{Tr} \int d^3x [\mathcal{G}(x, x, \omega + i\epsilon) - \mathcal{G}_0(x, x, \omega + i\epsilon)]$$

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$\mathcal{G}$  full Greens function for the fluctuating field

$\mathcal{G}_0$  free Greens function

Trace is over spin

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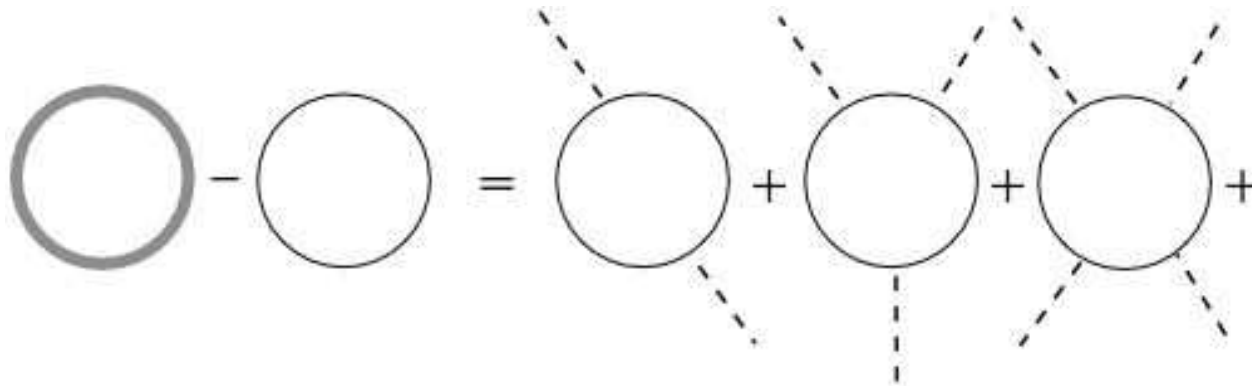
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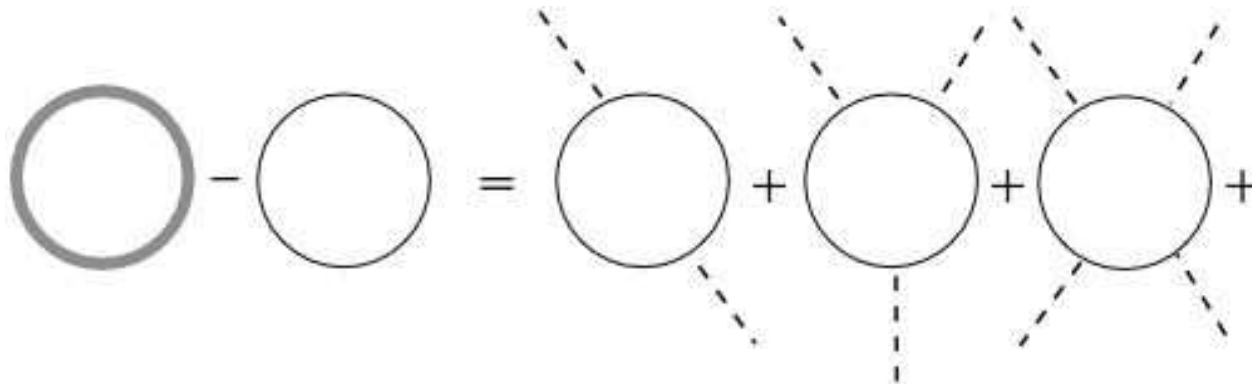
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⇒ “Experimental confirmation of the Casimir effect does not establish the reality of zero point fluctuations” [R. Jaffe et. al.]

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Moore; Razavy, Terning; Johnston, Sarkar; Dodonov et al;  
Plunien et al; Barton, Eberlein, Calogeracos; Ford, Vilenkin;  
Jaeckel, Reynaud, Lambrecht; Brevik, Milton et al;  
Dalvit, Maia-Neto et al; Law; Parentani, ...



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- The **dissipative** part we obtain **agrees** with the other methods. But those have problems with the **reactive** part, which in general yields a **non-positive** energy  $\implies$  **EXPERIMENT**

## SOME DETAILS OF THE METHOD

- **Hamiltonian method** for neutral Klein-Gordon field in a cavity  $\Omega_t$ , with boundaries moving at a certain speed  $v \ll c$ ,  $\epsilon = v/c$  (of order  $10^{-8}$  in **Kim, Brownell, Onofrio, PRL 96 (2006) 200402**)

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- **Hamiltonian.** Transform moving boundary into fixed one by (non-conformal) change of coordinates

$$\mathcal{R} : (\bar{t}, \mathbf{y}) \rightarrow (t(\bar{t}, \mathbf{y}), \mathbf{x}(\bar{t}, \mathbf{y})) = (\bar{t}, \mathbf{R}(\bar{t}, \mathbf{y}))$$

transform  $\Omega_t$  into a fixed domain  $\tilde{\Omega}$

$$\tilde{\Omega}: (t(\bar{t}, \mathbf{y}), \mathbf{x}(\bar{t}, \mathbf{y})) = \mathcal{R}(\bar{t}, \mathbf{y}) = (\bar{t}, \mathbf{R}(\bar{t}, \mathbf{y}))$$

(with  $\bar{t}$  the new time)

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renormalized energy **negative** while the mirror moves:

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Trajectory  $(t, \epsilon g(t))$ . When mirror at rest, **scattering** described by **matrix**

$$S(\omega) = \begin{pmatrix} s(\omega) & r(\omega) e^{-2i\omega L} \\ r(\omega) e^{2i\omega L} & s(\omega) \end{pmatrix}$$

$\implies S$  matrix is taken to be:  $(x = L$  position of the mirror)

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Trajectory  $(t, \epsilon g(t))$ . When mirror at rest, **scattering** described by **matrix**

$$S(\omega) = \begin{pmatrix} s(\omega) & r(\omega) e^{-2i\omega L} \\ r(\omega) e^{2i\omega L} & s(\omega) \end{pmatrix}$$

$\implies$   $S$  matrix is taken to be:  $(x = L$  position of the mirror)

$\rightarrow$  Real in the temporal domain:  $S(-\omega) = S^*(\omega)$

$\rightarrow$  Causal:  $S(\omega)$  is analytic for  $\text{Im}(\omega) > 0$

$\rightarrow$  Unitary:  $S(\omega)S^\dagger(\omega) = \text{Id}$

$\rightarrow$  The identity at high frequencies:  $S(\omega) \rightarrow \text{Id}$ , when  $|\omega| \rightarrow \infty$

$s(\omega)$  and  $r(\omega)$  **meromorphic** (cut-off) functions

(material's **permittivity** and **resistivity**)



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Note this integral **diverges** for a perfect mirror ( $r \equiv -1$ ,  $s \equiv 0$ , ideal case), but **nicely converges** for our partially transmitting (physical) one where  $r(\omega) \rightarrow 0$ ,  $s(\omega) \rightarrow 1$ , as  $\omega \rightarrow \infty$

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$\Rightarrow$  **Two** mirrors; **higher** dimensions; fields of **any** kind

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● The main issue: [S.A. Fulling et. al., hep-th/070209v2](#)

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- **Idea**: zero point fluctuations can contribute to the **cosmological constant** Ya.B. Zeldovich '68

## CC PROBLEM

- Relativistic field: collection of harmonic oscill's (scalar field)

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- What we **do consider** —with relative success in some different approaches— is the **additional** contribution to the cc coming from the **non-trivial topology** of space or from specific **boundary conditions** imposed on braneworld models:

⇒ kind of cosmological Casimir effect

# Cosmolog Imprint of the Casimir Eff't?

- Assuming one will be able to prove (in the future) that the ground value of the cc is **zero** (as many had suspected until recently), we will be left with this **incremental value** coming from the topology or BCs
  - \* **L. Parker & A. Raval**,  $\Lambda$ CDM, vacuum energy density
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  - **(c)** supergraviton theories (discret dims, deconstr)

# The Braneworld Case

1. Braneworld may help to solve:

- the hierarchy problem
- the cosmological constant problem

2. Presumably, the bulk Casimir effect will play a role in the construction (radion stabilization) of braneworlds

- Bulk Casimir effect (effective potential) for a conformal or massive scalar field
- Bulk is a 5-dim AdS or dS space with 2/1 4-dim dS brane (our universe)
- Consistent with observational data even for relatively large extra dimension

Previous work:

- flat space brane
- bulk conformal scalar field
- conclusion: no CE

We used **zeta regularization** at full power, with **positive** results!

EE, S Nojiri, SD Odintsov, S Ogushi, Phys Rev D67 (2003) 063515 *Casimir effect in de Sitter and Anti-de Sitter braneworlds* EE, SD Odintsov, AA Saharian 0902.0717

*Repulsive Casimir effect from extra dimensions and Robin BC: from branes to pistons*

# Gravity Eqs as Eqs of State: $f(R)$ Case

- The cosmological constant as an “integration constant”

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- C. Eling, R. Guedens, T. Jacobson [PRL2006]: extension to polynomial  $f(R)$  gravity but as non-equilibrium thermodyn.  
Also Erik Verlinde (private discussions)

● **Jacobson's argument:** basic thermodynamic relation

$$\delta Q = T\delta S$$

- entropy proportional to variation of the horizon area:  $\delta S = \eta \delta \mathcal{A}$
- local temperature  $T$  defined as **Unruh temp**:  $T = \hbar k / 2\pi$
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- **Case of  $\mathbf{f}(R)$  gravities:**  $\mathbf{L} = \mathbf{f}(R, \nabla^n R)$

- Also the concept of an **effective Newton constant** for graviton exchange (**effective propagator**)

$$\begin{aligned}\frac{1}{8\pi G_{eff}} &= E_R^{pqrs} \epsilon_{pq} \epsilon_{rs} = \frac{\partial \mathbf{f}}{\partial R} (g^{pr} g^{qs} - g^{qr} g^{ps}) \epsilon_{pq} \epsilon_{rs} \\ &= \frac{\partial \mathbf{f}}{\partial R} = \frac{\eta_e}{2\pi}, \quad S = \frac{A}{4 G_{eff}}\end{aligned}$$

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- For these theories, the different polarizations of the gravitons only enter in the definition of the **effective Newton constant through the metric itself**
- Final result, for  $\mathbf{f}(R)$  gravities:  
*the local field equations can be thought of as an equation of state of equilibrium thermodynamics* (as in the GR case)



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T Zhu, Ji-R Ren and S-F Mo, arXiv:0805.1162 [gr-qc];  
C Eling, arXiv:0806.3165 [hep-th]; R-G Cai, L-M Cao and Y-P Hu, arXiv:0807.1232 [hep-th] & arXiv:0809.1554 [hep-th]

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Happy Birthday, Iver! Long and Pleasant Life!