

ICE



IEEC

Vacuum Fluctuations: Regularization & Cosmological Issues

EMILIO ELIZALDE

ICE/CSIC & IEEC, UAB, Barcelona

ESF CASIMIR Topical Meeting, Wien, 11. Mai 2011

Outline

- On Einstein's Cosmological Constant: a Historical Perspective
- Quantum Vacuum Fluctuations: the Casimir Effect
- Vacuum Fluctuations and the Equivalence Principle
- The Sign of the Vacuum Forces
- Repulsion from Higher Dimensions and BCs
- CE and Accelerated Expansion (Dark Energy): a Cosmo-Topological Casimir Effect?
- Gravity Equations as Equations of State

Our Universe: brief historical account

- The present description of our Universe started to take form during the 3rd decade of last Century [E Mach, F Wilczek]

Our Universe: brief historical account

- The present description of our Universe started to take form during the 3rd decade of last Century [E Mach, F Wilczek]

- Einstein

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu}$$

Our Universe: brief historical account

- The present description of our Universe started to take form during the 3rd decade of last Century [E Mach, F Wilczek]

- Einstein

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu}$$

- But... the Universe was static !!

Our Universe: brief historical account

- The present description of our Universe started to take form during the 3rd decade of last Century [E Mach, F Wilczek]

- Einstein
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu}$$

- But... the Universe was static !!

- The cc λ
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu} + \lambda g_{\mu\nu}$$

Our Universe: brief historical account

- The present description of our Universe started to take form during the 3rd decade of last Century [E Mach, F Wilczek]

- Einstein
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu}$$

- But... the Universe was static !!

- The cc λ
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu} + \lambda g_{\mu\nu}$$

- What is λ ? Non-physical

Our Universe: brief historical account

- The present description of our Universe started to take form during the 3rd decade of last Century [E Mach, F Wilczek]

- Einstein
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu}$$

- But... the Universe was static !!

- The cc λ
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu} + \lambda g_{\mu\nu}$$

- What is λ ? Non-physical

- Karl Schwarzschild: Black Hole solution (22 December 1915)

Our Universe: brief historical account

- The present description of our Universe started to take form during the 3rd decade of last Century [E Mach, F Wilczek]

- Einstein
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu}$$

- But... the Universe was static !!

- The cc λ
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu} + \lambda g_{\mu\nu}$$

- What is λ ? Non-physical

- Karl Schwarzschild: Black Hole solution (22 December 1915)

- Alexander Friedmann: expanding universe solution (1922)

Our Universe: brief historical account

- The present description of our Universe started to take form during the 3rd decade of last Century [E Mach, F Wilczek]

- Einstein
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu}$$

- But... the Universe was static !!

- The cc λ
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu} + \lambda g_{\mu\nu}$$

- What is λ ? Non-physical

- Karl Schwarzschild: Black Hole solution (22 December 1915)

- Alexander Friedmann: expanding universe solution (1922)

- Willem de Sitter: massless universe static solution (just cc) 'dark matter' (with AE, 1932)

Our Universe: brief historical account

- The present description of our Universe started to take form during the 3rd decade of last Century [E Mach, F Wilczek]

- Einstein
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu}$$

- But... the Universe was static !!

- The cc λ
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu} + \lambda g_{\mu\nu}$$

- What is λ ? Non-physical

- Karl Schwarzschild: Black Hole solution (22 December 1915)

- Alexander Friedmann: expanding universe solution (1922)

- Willem de Sitter: massless universe static solution (just cc) 'dark matter' (with AE, 1932)

- Georges Lemaître: expanding universe (MIT 1925, AF sol); visited Vesto Slipher (Lowell Obs, Arizona, 1912 galaxy redshifts) and Edwin Hubble (Mount Wilson, Pasadena); Keeler-Slipher-Campbell, 1918

Our Universe: brief historical account

- The present description of our Universe started to take form during the 3rd decade of last Century [E Mach, F Wilczek]

- Einstein
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu}$$

- But... the Universe was static !!

- The cc λ
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -(8\pi G/c^4)T_{\mu\nu} + \lambda g_{\mu\nu}$$

- What is λ ? Non-physical

- Karl Schwarzschild: Black Hole solution (22 December 1915)

- Alexander Friedmann: expanding universe solution (1922)

- Willem de Sitter: massless universe static solution (just cc) 'dark matter' (with AE, 1932)

- Georges Lemaître: expanding universe (MIT 1925, AF sol); visited Vesto Slipher (Lowell Obs, Arizona, 1912 galaxy redshifts) and Edwin Hubble (Mount Wilson, Pasadena); Keeler-Slipher-Campbell, 1918

- Led to Big Bang theory (Fred Hoyle, BBC radio's Third Programme, 18:30 GMT, 28 March 1949)

Historical account continued

- Fred Hoyle, Thomas Gold, Hermann Bondi: **Steady State** theory '48, "C-field" with **negative pressure** to be consistent with conservation of energy (anticipated inflation)

Historical account continued

- Fred Hoyle, Thomas Gold, Hermann Bondi: **Steady State** theory '48, "C-field" with **negative pressure** to be consistent with conservation of energy (anticipated inflation)
- **Friedmann-Lemaître-Robertson-Walker: FLRW universe** (1931,35-37)

Historical account continued

- Fred Hoyle, Thomas Gold, Hermann Bondi: **Steady State** theory '48, "C-field" with **negative pressure** to be consistent with conservation of energy (anticipated inflation)
- Friedmann-Lemaître-Robertson-Walker: **FLRW universe** (1931,35-37)
- **Dark matter**: postulated by **Fritz Zwicky** (1934) to account for evidence of 'missing mass' in the orbital velocities of galaxies in clusters. Now lots of evidence: galactic rotation curves (**Vera Rubin, Kent Ford**, 1975), gravitational lensing, etc.

Historical account continued

- Fred Hoyle, Thomas Gold, Hermann Bondi: **Steady State** theory '48, "C-field" with **negative pressure** to be consistent with conservation of energy (anticipated inflation)
- **Friedmann-Lemaître-Robertson-Walker: FLRW universe** (1931,35-37)
- **Dark matter:** postulated by **Fritz Zwicky** (1934) to account for evidence of 'missing mass' in the orbital velocities of galaxies in clusters. Now lots of evidence: galactic rotation curves (**Vera Rubin, Kent Ford**, 1975), gravitational lensing, etc.
- **Arno Penzias, Robert Wilson** (1964 Bell Labs, New Jersey); at Princeton (< 40 miles away) **Dicke, Peebles, Wilkinson** writing paper on how CMB should be
Burke (MIT) told Penzias of Peebles' work: and **the Big Bang was here !**

Historical account continued

- Fred Hoyle, Thomas Gold, Hermann Bondi: **Steady State** theory '48, "C-field" with **negative pressure** to be consistent with conservation of energy (anticipated inflation)
- **Friedmann-Lemaître-Robertson-Walker: FLRW universe** (1931,35-37)
- **Dark matter:** postulated by **Fritz Zwicky** (1934) to account for evidence of 'missing mass' in the orbital velocities of galaxies in clusters. Now lots of evidence: galactic rotation curves (**Vera Rubin, Kent Ford**, 1975), gravitational lensing, etc.
- **Arno Penzias, Robert Wilson** (1964 Bell Labs, New Jersey); at Princeton (< 40 miles away) **Dicke, Peebles, Wilkinson** writing paper on how CMB should be **Burke (MIT)** told Penzias of Peebles' work: and **the Big Bang was here !**
- **Alan Guth** (January 1980), **Andrei Linde: inflation.** Also: **Alexei Starobinsky, Andreas Albrecht, Paul Steinhardt.** And **Zel'dovich, Coleman** (decay of **false vacuum**)

Historical account continued

- Fred Hoyle, Thomas Gold, Hermann Bondi: **Steady State** theory '48, "C-field" with **negative pressure** to be consistent with conservation of energy (anticipated inflation)
- **Friedmann-Lemaître-Robertson-Walker: FLRW universe** (1931,35-37)
- **Dark matter**: postulated by **Fritz Zwicky** (1934) to account for evidence of 'missing mass' in the orbital velocities of galaxies in clusters. Now lots of evidence: galactic rotation curves (**Vera Rubin, Kent Ford**, 1975), gravitational lensing, etc.
- **Arno Penzias, Robert Wilson** (1964 Bell Labs, New Jersey); at Princeton (< 40 miles away) **Dicke, Peebles, Wilkinson** writing paper on how CMB should be **Burke (MIT)** told Penzias of Peebles' work: and **the Big Bang was here !**
- **Alan Guth** (January 1980), **Andrei Linde: inflation**. Also: **Alexei Starobinsky, Andreas Albrecht, Paul Steinhardt**. And **Zel'dovich, Coleman** (decay of **false vacuum**)
- Quantum fluctuations + inflation \longrightarrow **multiverse**
Hawking+Turok instanton '98, **NoBoundary (HH)** vs **chaotic inflation**

Historical account continued

- Fred Hoyle, Thomas Gold, Hermann Bondi: **Steady State** theory '48, "C-field" with **negative pressure** to be consistent with conservation of energy (anticipated inflation)
- **Friedmann-Lemaître-Robertson-Walker: FLRW universe** (1931,35-37)
- **Dark matter**: postulated by **Fritz Zwicky** (1934) to account for evidence of 'missing mass' in the orbital velocities of galaxies in clusters. Now lots of evidence: galactic rotation curves (**Vera Rubin, Kent Ford**, 1975), gravitational lensing, etc.
- **Arno Penzias, Robert Wilson** (1964 Bell Labs, New Jersey); at Princeton (< 40 miles away) **Dicke, Peebles, Wilkinson** writing paper on how CMB should be **Burke (MIT)** told Penzias of Peebles' work: and **the Big Bang was here !**
- **Alan Guth** (January 1980), **Andrei Linde: inflation**. Also: **Alexei Starobinsky, Andreas Albrecht, Paul Steinhardt**. And **Zel'dovich, Coleman** (decay of **false vacuum**)
- Quantum fluctuations + inflation \longrightarrow **multiverse**
Hawking+Turok instanton '98, **NoBoundary** (HH) vs **chaotic inflation**
- **Perlmutter ea, Riess ea**, 1998: Universe expansion is **accelerating !**

Historical account continued

- Fred Hoyle, Thomas Gold, Hermann Bondi: **Steady State** theory '48, "C-field" with **negative pressure** to be consistent with conservation of energy (anticipated inflation)
- Friedmann-Lemaître-Robertson-Walker: **FLRW universe** (1931,35-37)
- **Dark matter**: postulated by **Fritz Zwicky** (1934) to account for evidence of 'missing mass' in the orbital velocities of galaxies in clusters. Now lots of evidence: galactic rotation curves (**Vera Rubin, Kent Ford**, 1975), gravitational lensing, etc.
- **Arno Penzias, Robert Wilson** (1964 Bell Labs, New Jersey); at Princeton (< 40 miles away) **Dicke, Peebles, Wilkinson** writing paper on how CMB should be **Burke (MIT)** told Penzias of Peebles' work: and **the Big Bang was here !**
- **Alan Guth** (January 1980), **Andrei Linde**: **inflation**. Also: **Alexei Starobinsky, Andreas Albrecht, Paul Steinhardt**. And **Zel'dovich, Coleman** (decay of **false vacuum**)
- Quantum fluctuations + inflation \longrightarrow **multiverse**
Hawking+Turok instanton '98, **NoBoundary** (HH) vs **chaotic inflation**
- **Perlmutter ea, Riess ea**, 1998: Universe expansion is **accelerating !**
- **Dark Energy**: scalar-tensor, $f(R)$ theories // **COBE, WMAP, PLANCK (LISA, BBO,...)**

Historical account continued

- Fred Hoyle, Thomas Gold, Hermann Bondi: **Steady State** theory '48, "C-field" with **negative pressure** to be consistent with conservation of energy (anticipated inflation)
- **Friedmann-Lemaître-Robertson-Walker: FLRW universe** (1931,35-37)
- **Dark matter:** postulated by **Fritz Zwicky** (1934) to account for evidence of 'missing mass' in the orbital velocities of galaxies in clusters. Now lots of evidence: galactic rotation curves (**Vera Rubin, Kent Ford**, 1975), gravitational lensing, etc.
- **Arno Penzias, Robert Wilson** (1964 Bell Labs, New Jersey); at Princeton (< 40 miles away) **Dicke, Peebles, Wilkinson** writing paper on how CMB should be **Burke (MIT)** told Penzias of Peebles' work: and **the Big Bang was here !**
- **Alan Guth** (January 1980), **Andrei Linde: inflation**. Also: **Alexei Starobinsky, Andreas Albrecht, Paul Steinhardt**. And **Zel'dovich, Coleman** (decay of **false vacuum**)
- Quantum fluctuations + inflation \longrightarrow **multiverse**
Hawking+Turok instanton '98, **NoBoundary** (HH) vs **chaotic inflation**
- **Perlmutter ea, Riess ea**, 1998: Universe expansion is **accelerating !**
- **Dark Energy:** scalar-tensor, $f(R)$ theories // **COBE, WMAP, PLANCK (LISA, BBO,...)**
- On a different level: Richard Dawkins, Hoyle's fallacy
evolutionary biology \longleftrightarrow **intelligent design**

Trying to solve these puzzles !

- The cc λ is indeed a peculiar quantity

Trying to solve these puzzles !

- The cc λ is indeed a peculiar quantity
- has to do with cosmology Einstein's eqs., FRW universe

Trying to solve these puzzles !

- The cc λ is indeed a peculiar quantity
 - has to do with cosmology Einstein's eqs., FRW universe
 - has to do with the local structure of elementary particle physics
stress-energy density μ of the vacuum

$$L_{cc} = \int d^4x \sqrt{-g} \mu^4 = \frac{1}{8\pi G} \int d^4x \sqrt{-g} \lambda$$

Trying to solve these puzzles !

- The cc λ is indeed a peculiar quantity
 - has to do with cosmology Einstein's eqs., FRW universe
 - has to do with the local structure of elementary particle physics stress-energy density μ of the vacuum

$$L_{cc} = \int d^4x \sqrt{-g} \mu^4 = \frac{1}{8\pi G} \int d^4x \sqrt{-g} \lambda$$

- In other words: two contributions on the same footing [Pauli 20's, Zel'dovich '68]

$$\frac{\Lambda c^2}{8\pi G} + \frac{1}{\text{Vol}} \frac{\hbar c}{2} \sum_i \omega_i$$

Trying to solve these puzzles !

- The cc λ is indeed a peculiar quantity
 - has to do with cosmology Einstein's eqs., FRW universe
 - has to do with the local structure of elementary particle physics stress-energy density μ of the vacuum

$$L_{cc} = \int d^4x \sqrt{-g} \mu^4 = \frac{1}{8\pi G} \int d^4x \sqrt{-g} \lambda$$

- In other words: two contributions on the same footing [Pauli 20's, Zel'dovich '68]

$$\frac{\Lambda c^2}{8\pi G} + \frac{1}{\text{Vol}} \frac{\hbar c}{2} \sum_i \omega_i$$

- For elementary particle physicists: a great embarrassment
no way to get rid off
Coleman, Hawking, Weinberg, Polchinski, ... '88-'89

Trying to solve these puzzles !

- The cc λ is indeed a peculiar quantity
 - has to do with cosmology Einstein's eqs., FRW universe
 - has to do with the local structure of elementary particle physics stress-energy density μ of the vacuum

$$L_{cc} = \int d^4x \sqrt{-g} \mu^4 = \frac{1}{8\pi G} \int d^4x \sqrt{-g} \lambda$$

- In other words: two contributions on the same footing [Pauli 20's, Zel'dovich '68]

$$\frac{\Lambda c^2}{8\pi G} + \frac{1}{\text{Vol}} \frac{\hbar c}{2} \sum_i \omega_i$$

- For elementary particle physicists: a great embarrassment
no way to get rid off
Coleman, Hawking, Weinberg, Polchinski, ... '88-'89

THE COSMOLOGICAL CONSTANT PROBLEM

Zero point energy

QFT vacuum to vacuum transition: $\langle 0|H|0\rangle$

Zero point energy

QFT vacuum to vacuum transition: $\langle 0|H|0\rangle$

Spectrum, normal ordering (harm oscill):

$$H = \left(n + \frac{1}{2} \right) \lambda_n a_n a_n^\dagger$$

Zero point energy

QFT vacuum to vacuum transition: $\langle 0|H|0\rangle$

Spectrum, normal ordering (harm oscill):

$$H = \left(n + \frac{1}{2} \right) \lambda_n a_n a_n^\dagger$$

$$\langle 0|H|0\rangle = \frac{\hbar c}{2} \sum_n \lambda_n = \frac{1}{2} \text{tr } H$$

Zero point energy

QFT vacuum to vacuum transition: $\langle 0|H|0\rangle$

Spectrum, normal ordering (harm oscill):

$$H = \left(n + \frac{1}{2} \right) \lambda_n a_n a_n^\dagger$$

$$\langle 0|H|0\rangle = \frac{\hbar c}{2} \sum_n \lambda_n = \frac{1}{2} \text{tr } H$$

gives ∞ physical meaning?

Zero point energy

QFT vacuum to vacuum transition: $\langle 0|H|0\rangle$

Spectrum, normal ordering (harm oscill):

$$H = \left(n + \frac{1}{2} \right) \lambda_n a_n a_n^\dagger$$

$$\langle 0|H|0\rangle = \frac{\hbar c}{2} \sum_n \lambda_n = \frac{1}{2} \text{tr } H$$

gives ∞ physical meaning?

Regularization + Renormalization (cut-off, dim, ζ)

Zero point energy

QFT vacuum to vacuum transition: $\langle 0|H|0\rangle$

Spectrum, normal ordering (harm oscill):

$$H = \left(n + \frac{1}{2} \right) \lambda_n a_n a_n^\dagger$$

$$\langle 0|H|0\rangle = \frac{\hbar c}{2} \sum_n \lambda_n = \frac{1}{2} \text{tr } H$$

gives ∞ physical meaning?

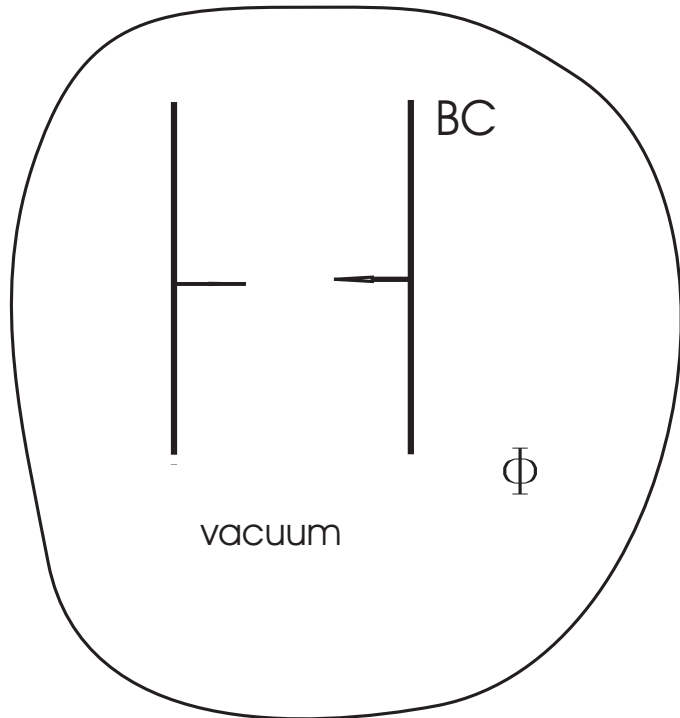
Regularization + Renormalization (cut-off, dim, ζ)

Even then: Has the final value real sense ?

The Casimir Effect

The Casimir Effect

BC e.g. periodic

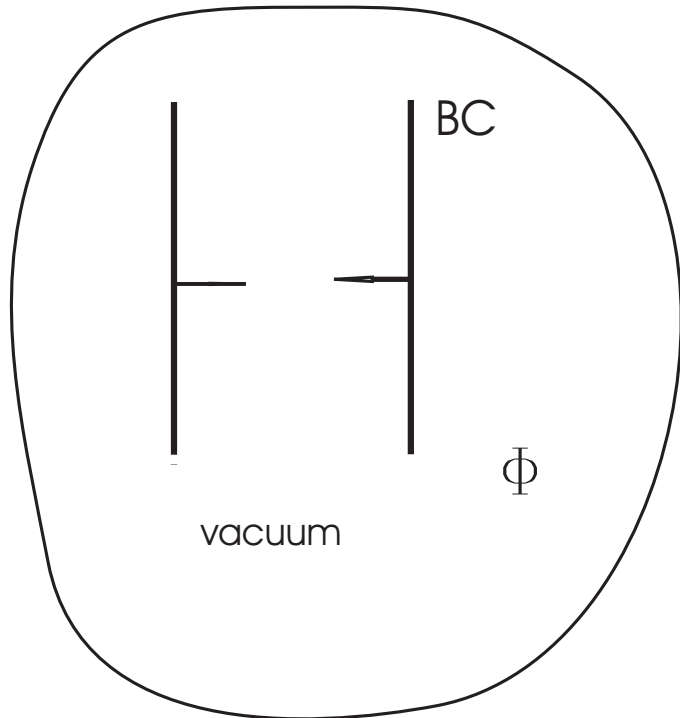


Casimir Effect

The Casimir Effect

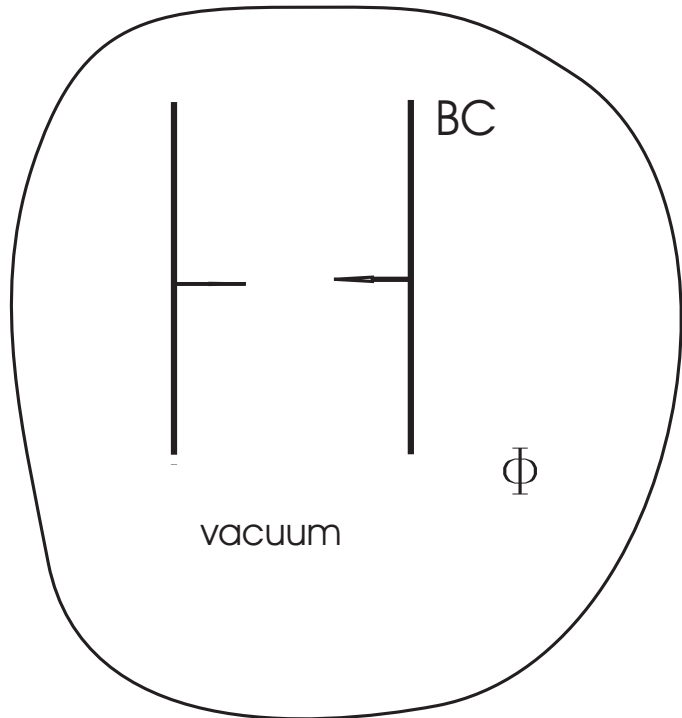
BC e.g. periodic

\Rightarrow all kind of fields



Casimir Effect

The Casimir Effect



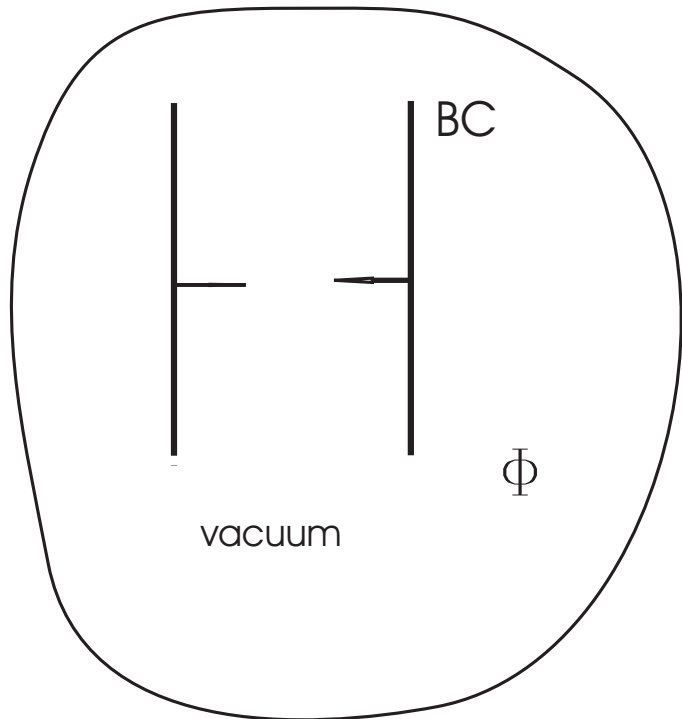
Casimir Effect

BC e.g. periodic

\Rightarrow all kind of fields

\Rightarrow curvature or topology

The Casimir Effect



Casimir Effect

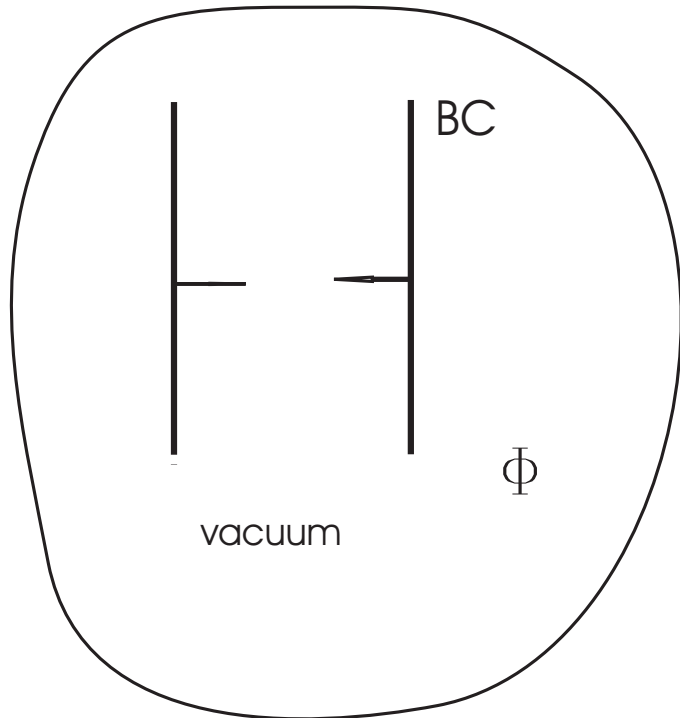
BC e.g. periodic

\Rightarrow all kind of fields

\Rightarrow curvature or topology

Universal process:

The Casimir Effect



Casimir Effect

BC e.g. periodic

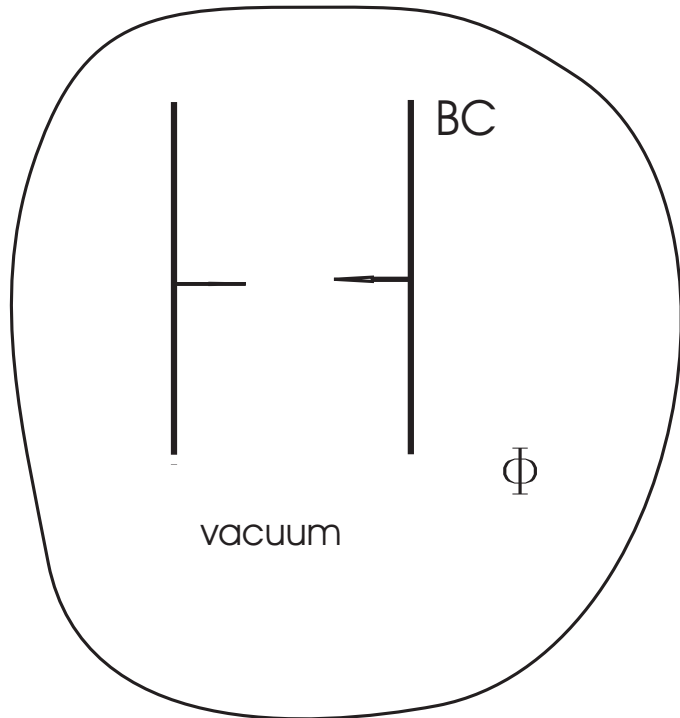
⇒ all kind of fields

⇒ curvature or topology

Universal process:

- Sonoluminescence (Schwinger)
- Cond. matter (wetting ^3He alc.)
- Optical cavities
- Direct experim. confirmation

The Casimir Effect



Casimir Effect

BC e.g. periodic

⇒ all kind of fields

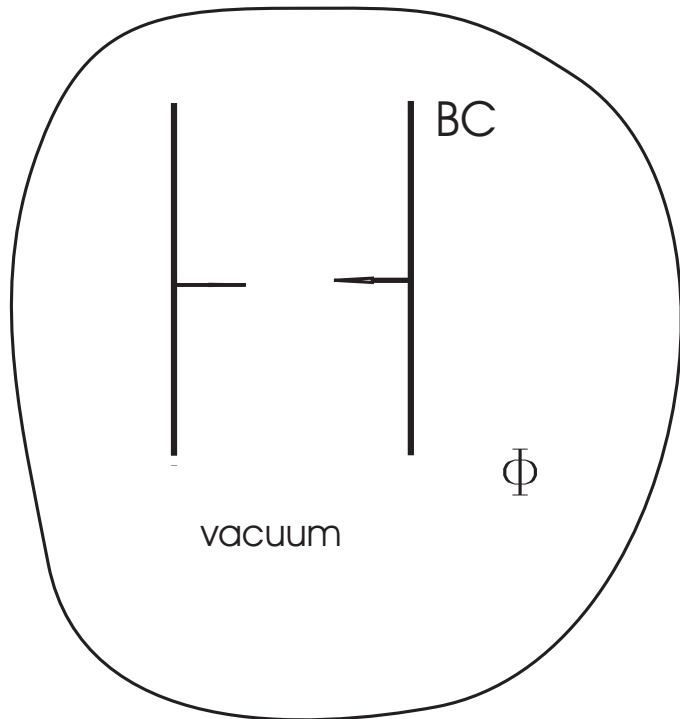
⇒ curvature or topology

Universal process:

- Sonoluminescence (Schwinger)
- Cond. matter (wetting ^3He alc.)
- Optical cavities
- Direct experim. confirmation

Van der Waals, Lifschitz theory

The Casimir Effect



Casimir Effect

- BC e.g. periodic
- \Rightarrow all kind of fields
- \Rightarrow curvature or topology

Universal process:

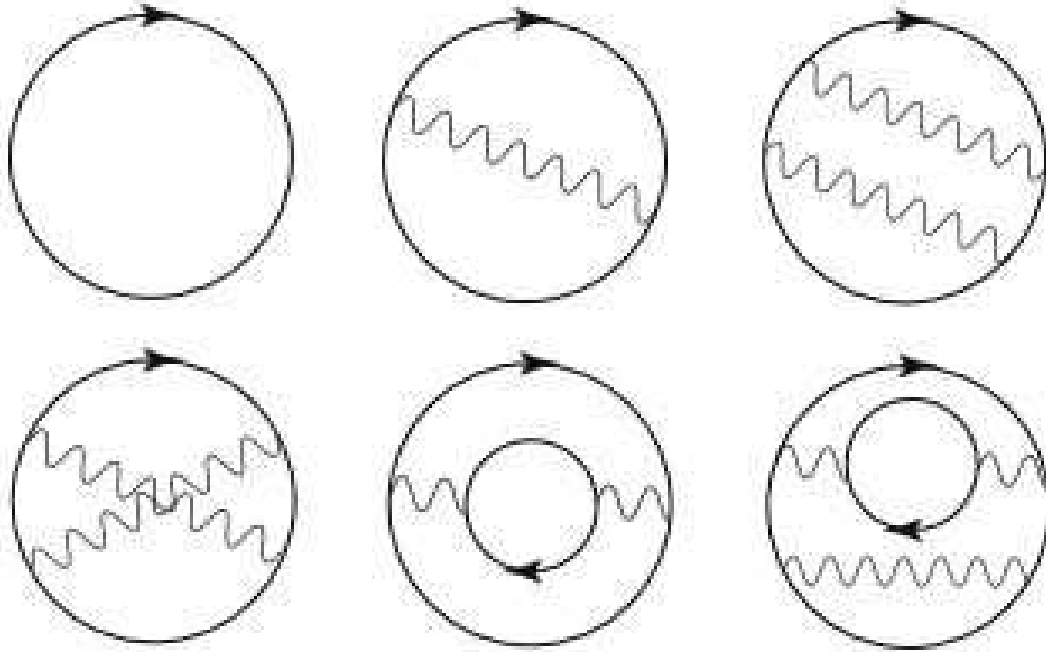
- Sonoluminescence (Schwinger)
- Cond. matter (wetting ^3He alc.)
- Optical cavities
- Direct experim. confirmation

Van der Waals, Lifschitz theory

- Dynamical CE \Leftarrow
- Lateral CE
- Extract energy from vacuum
- CE and the cosmological constant \Leftarrow

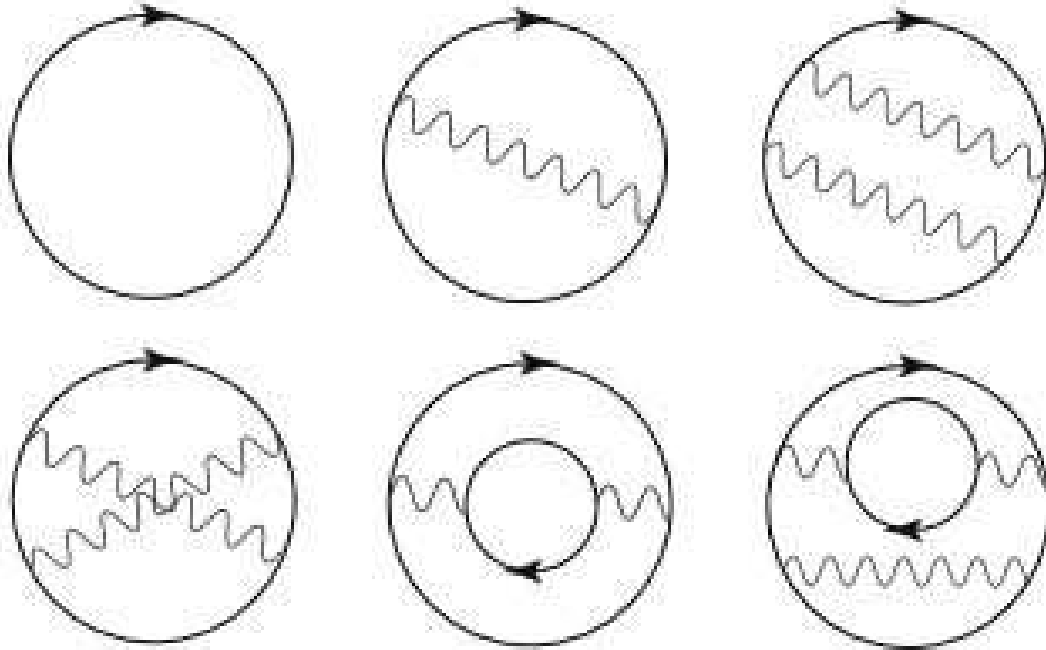
The standard approach

The standard approach



⇒ Casimir force: calculated by computing change in zero point energy of the em field

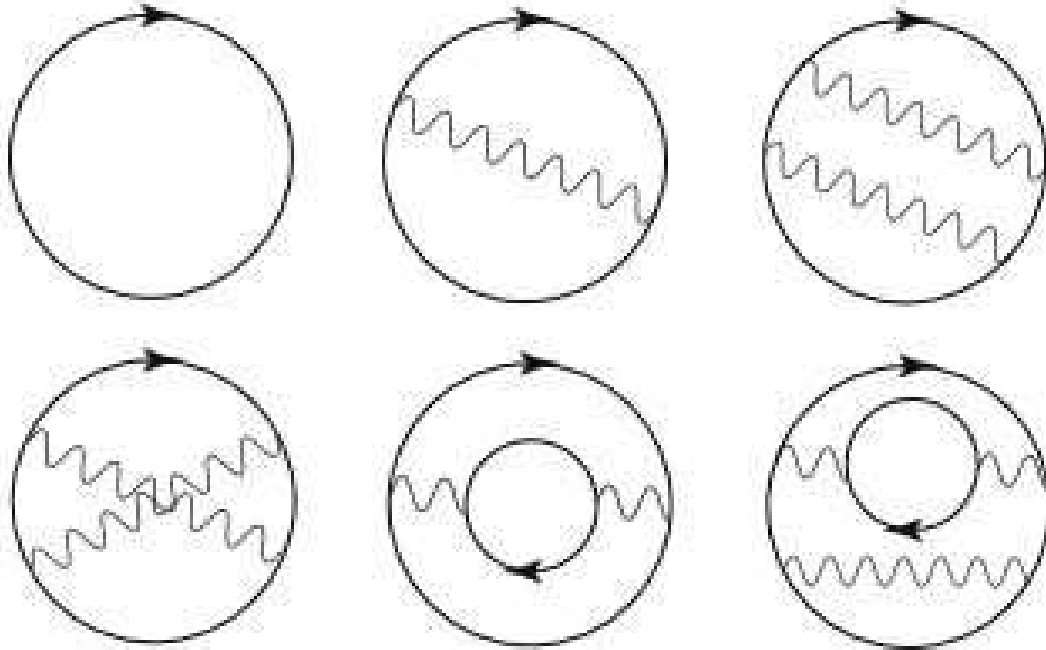
The standard approach



⇒ Casimir force: calculated by computing change in zero point energy of the em field

⇒ But Casimir effects can be calculated as S -matrix elements:
Feynman diagrs with ext. lines

The standard approach



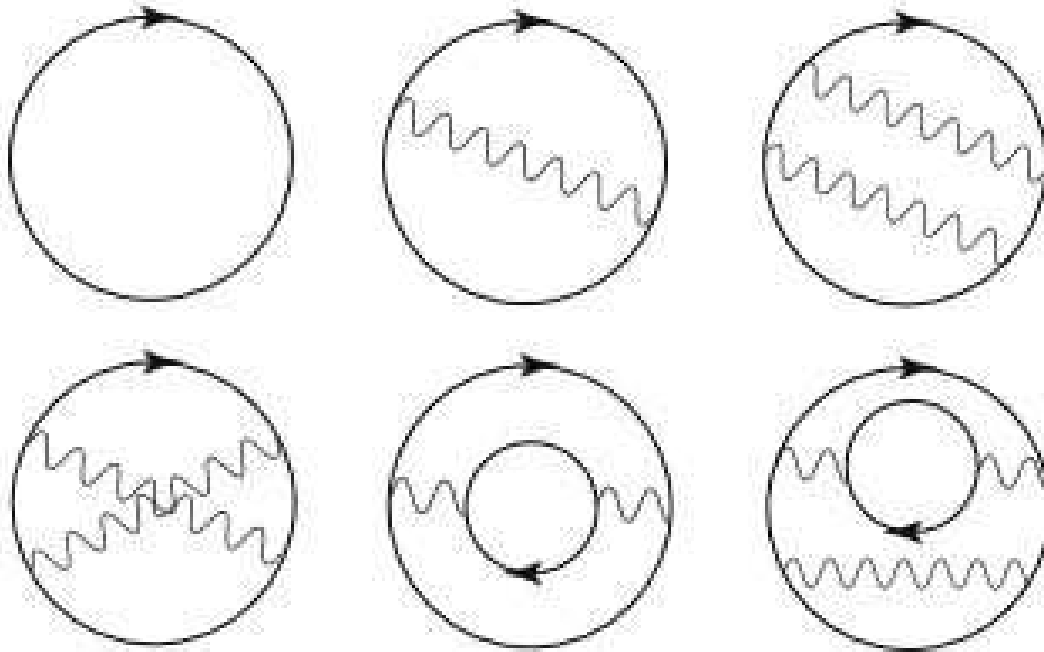
⇒ Casimir force: calculated by computing change in zero point energy of the em field

⇒ But Casimir effects can be calculated as S -matrix elements: Feynman diagrs with ext. lines

In modern language the Casimir energy can be expressed in terms of the trace of the Greens function for the fluctuating field in the background of interest (conducting plates)

$$\mathcal{E} = \frac{\hbar}{2\pi} \text{Im} \int d\omega \omega \text{Tr} \int d^3x [\mathcal{G}(x, x, \omega + i\epsilon) - \mathcal{G}_0(x, x, \omega + i\epsilon)]$$

The standard approach



⇒ Casimir force: calculated by computing change in zero point energy of the em field

⇒ But Casimir effects can be calculated as S -matrix elements: Feynman diagrs with ext. lines

In modern language the Casimir energy can be expressed in terms of the trace of the Greens function for the fluctuating field in the background of interest (conducting plates)

$$\mathcal{E} = \frac{\hbar}{2\pi} \text{Im} \int d\omega \omega \text{Tr} \int d^3x [\mathcal{G}(x, x, \omega + i\epsilon) - \mathcal{G}_0(x, x, \omega + i\epsilon)]$$

\mathcal{G} full Greens function for the fluctuating field

\mathcal{G}_0 free Greens function

Trace is over spin

$$E_C = \langle \quad \rangle_{\text{plates}} - \langle \quad \rangle_{\text{no plates}}$$

$$E_C = \langle \quad \rangle_{\text{plates}} - \langle \quad \rangle_{\text{no plates}}$$

$$\frac{1}{\pi} \text{Im} \int [\mathcal{G}(x, x, \omega + i\epsilon) - \mathcal{G}_0(x, x, \omega + i\epsilon)] = \frac{d\Delta N}{d\omega}$$

change in the density of states due to the background

$$E_C = \langle \quad \rangle_{\text{plates}} - \langle \quad \rangle_{\text{no plates}}$$

$$\frac{1}{\pi} \text{Im} \int [\mathcal{G}(x, x, \omega + i\epsilon) - \mathcal{G}_0(x, x, \omega + i\epsilon)] = \frac{d\Delta N}{d\omega}$$

change in the density of states due to the background

⇒ A restatement of the Casimir sum over shifts in zero-point energies

$$\frac{\hbar}{2} \sum (\omega - \omega_0)$$

$$E_C = \langle \quad \rangle_{\text{plates}} - \langle \quad \rangle_{\text{no plates}}$$

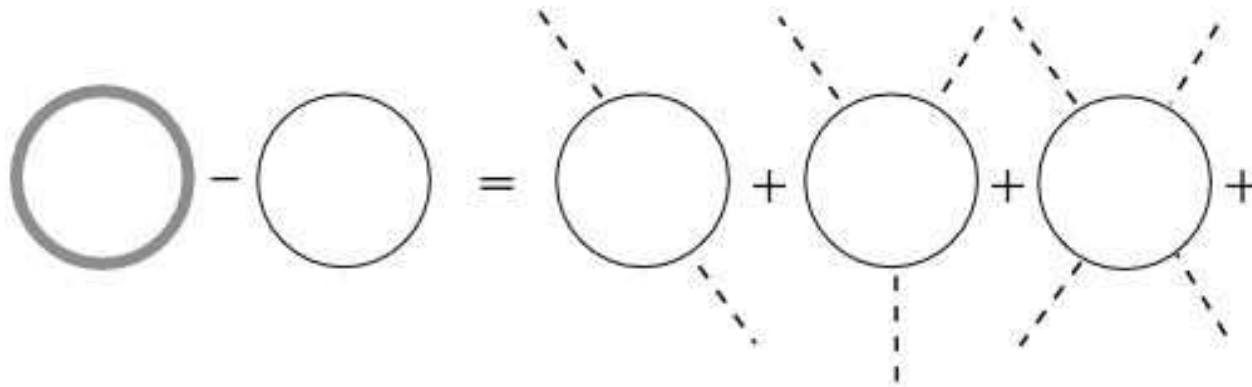
$$\frac{1}{\pi} \text{Im} \int [\mathcal{G}(x, x, \omega + i\epsilon) - \mathcal{G}_0(x, x, \omega + i\epsilon)] = \frac{d\Delta N}{d\omega}$$

change in the density of states due to the background

⇒ A restatement of the Casimir sum over shifts in zero-point energies

$$\frac{\hbar}{2} \sum (\omega - \omega_0)$$

⇒ **Lippman-Schwinger eq.** allows full Greens f, \mathcal{G} , be expanded as a series in free Green's f, \mathcal{G}_0 , and the coupling to the external field



$$E_C = \langle \quad \rangle_{\text{plates}} - \langle \quad \rangle_{\text{no plates}}$$

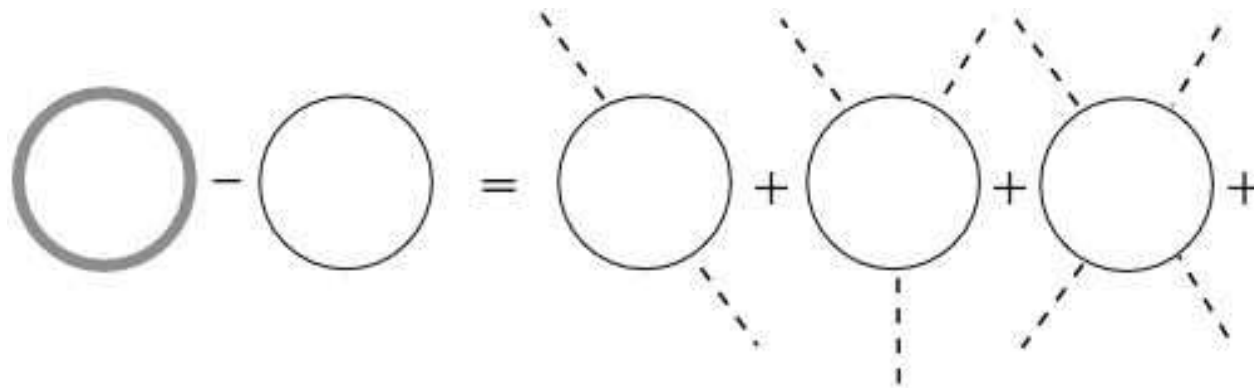
$$\frac{1}{\pi} \text{Im} \int [\mathcal{G}(x, x, \omega + i\epsilon) - \mathcal{G}_0(x, x, \omega + i\epsilon)] = \frac{d\Delta N}{d\omega}$$

change in the density of states due to the background

⇒ A restatement of the Casimir sum over shifts in zero-point energies

$$\frac{\hbar}{2} \sum (\omega - \omega_0)$$

⇒ **Lippman-Schwinger eq.** allows full Greens f, \mathcal{G} , be expanded as a series in free Green's f, \mathcal{G}_0 , and the coupling to the external field



⇒ “Experimental confirmation of the Casimir effect doesn’t establish the reality of zero point fluct’s better than say the Lamb shift does” [R Jaffe e a]

Vacuum Fluct & the Equival Principle

● The main issue:

S.A. Fulling et. al., hep-th/070209

energy **ALWAYS gravitates** therefore the energy density of the vacuum appears on the rhs of Einstein's equations:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -8\pi G(\tilde{T}_{\mu\nu} - \mathcal{E}g_{\mu\nu})$$

Vacuum Fluct & the Equival Principle

- The main issue:

S.A. Fulling et. al., hep-th/070209

energy **ALWAYS gravitates** therefore the energy density of the vacuum appears on the rhs of Einstein's equations:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -8\pi G(\tilde{T}_{\mu\nu} - \mathcal{E}g_{\mu\nu})$$

- Equivalent to a **cosmological const** $\Lambda = 8\pi G\mathcal{E}$, $\rho_c = \frac{3H^2}{8\pi G}$

Vacuum Fluct & the Equival Principle

- The main issue: [S.A. Fulling et. al., hep-th/070209](#)

energy **ALWAYS gravitates** therefore the energy density of the vacuum appears on the rhs of Einstein's equations:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -8\pi G(\tilde{T}_{\mu\nu} - \mathcal{E}g_{\mu\nu})$$

- Equivalent to a **cosmological const** $\Lambda = 8\pi G\mathcal{E}$, $\rho_c = \frac{3H^2}{8\pi G}$

- **Observations:** [M. Tegmark et al. \[SDSS Collab.\] PRD 2004](#)

$$\Lambda = (2.14 \pm 0.13 \times 10^{-3} \text{ eV})^4 \sim 4.32 \times 10^{-9} \text{ erg/cm}^3$$

Vacuum Fluct & the Equival Principle

- The main issue: [S.A. Fulling et. al., hep-th/070209](#)

energy **ALWAYS gravitates** therefore the energy density of the vacuum appears on the rhs of Einstein's equations:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -8\pi G(\tilde{T}_{\mu\nu} - \mathcal{E}g_{\mu\nu})$$

- Equivalent to a **cosmological const** $\Lambda = 8\pi G\mathcal{E}$, $\rho_c = \frac{3H^2}{8\pi G}$

- **Observations:** [M. Tegmark et al. \[SDSS Collab.\] PRD 2004](#)

$$\Lambda = (2.14 \pm 0.13 \times 10^{-3} \text{ eV})^4 \sim 4.32 \times 10^{-9} \text{ erg/cm}^3$$

- **Question:** how finite Casimir energy of pair of plates **couple**s to gravity?

Vacuum Fluct & the Equival Principle

- The main issue: [S.A. Fulling et. al., hep-th/070209](#)

energy **ALWAYS gravitates** therefore the energy density of the vacuum appears on the rhs of Einstein's equations:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -8\pi G(\tilde{T}_{\mu\nu} - \mathcal{E}g_{\mu\nu})$$

- Equivalent to a **cosmological const** $\Lambda = 8\pi G\mathcal{E}$, $\rho_c = \frac{3H^2}{8\pi G}$

- **Observations:** [M. Tegmark et al. \[SDSS Collab.\] PRD 2004](#)

$$\Lambda = (2.14 \pm 0.13 \times 10^{-3} \text{ eV})^4 \sim 4.32 \times 10^{-9} \text{ erg/cm}^3$$

- **Question:** how finite Casimir energy of pair of plates **couple**s to gravity?

- **Two ways** to proceed. **Gauge-invariant** procedure:

energy-momentum tensor of the phys sys must be conserved, so include a physical mechanism holding the plates apart against the Casimir force

→ Leads to **complicated** model-dependent calculations

Vacuum Fluct & the Equival Principle

- The main issue: [S.A. Fulling et. al., hep-th/070209](#)

energy **ALWAYS gravitates** therefore the energy density of the vacuum appears on the rhs of Einstein's equations:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -8\pi G(\tilde{T}_{\mu\nu} - \mathcal{E}g_{\mu\nu})$$

- Equivalent to a **cosmological const** $\Lambda = 8\pi G\mathcal{E}$, $\rho_c = \frac{3H^2}{8\pi G}$

- **Observations:** [M. Tegmark et al. \[SDSS Collab.\] PRD 2004](#)

$$\Lambda = (2.14 \pm 0.13 \times 10^{-3} \text{ eV})^4 \sim 4.32 \times 10^{-9} \text{ erg/cm}^3$$

- **Question:** how finite Casimir energy of pair of plates **couples** to gravity?

- **Two ways** to proceed. **Gauge-invariant** procedure:

energy-momentum tensor of the phys sys must be conserved, so include a physical mechanism holding the plates apart against the Casimir force

→ Leads to **complicated** model-dependent calculations

- Alternative: find a **physically natural** coordinate system, more realistic than another: **Fermi** coord system [[Marzlin '94](#)]

Vacuum Fluct & the Equival Principle

- The main issue: [S.A. Fulling et. al., hep-th/070209](#)

energy **ALWAYS gravitates** therefore the energy density of the vacuum appears on the rhs of Einstein's equations:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -8\pi G(\tilde{T}_{\mu\nu} - \mathcal{E}g_{\mu\nu})$$

- Equivalent to a **cosmological const** $\Lambda = 8\pi G\mathcal{E}$, $\rho_c = \frac{3H^2}{8\pi G}$

- **Observations:** [M. Tegmark et al. \[SDSS Collab.\] PRD 2004](#)

$$\Lambda = (2.14 \pm 0.13 \times 10^{-3} \text{ eV})^4 \sim 4.32 \times 10^{-9} \text{ erg/cm}^3$$

- **Question:** how finite Casimir energy of pair of plates **couple** to gravity?

- **Two ways** to proceed. **Gauge-invariant** procedure:

energy-momentum tensor of the phys sys must be conserved, so include a physical mechanism holding the plates apart against the Casimir force

→ Leads to **complicated** model-dependent calculations

- Alternative: find a **physically natural** coordinate system, more realistic than another: **Fermi** coord system [[Marzlin '94](#)]

- Calculations done also in **Rindler coord** (uniform accel obs)

CC PROBLEM

- Relativistic field: collection of harmonic oscill's (scalar field)

$$E_0 = \frac{\hbar c}{2} \sum_n \omega_n, \quad \omega = k^2 + m^2/\hbar^2, \quad k = 2\pi/\lambda$$

CC PROBLEM

- Relativistic field: collection of harmonic oscill's (scalar field)

$$E_0 = \frac{\hbar c}{2} \sum_n \omega_n, \quad \omega = k^2 + m^2/\hbar^2, \quad k = 2\pi/\lambda$$

- Evaluating in a box and putting a cut-off at maximum k_{max} corresp'ng to QFT physics (e.g., Planck energy)

$$M_P/M_{ew} \sim 10^{16}, \quad M_P/M_{cc} \sim 10^{31}, \quad \rho \sim \frac{\hbar k_{Planck}^4}{16\pi^2} \sim 10^{123} \rho_{obs}$$

a thick **aether!** **Caldwell, Carroll** but **Gómez, Dvali**: species $\downarrow 10^{30}$

CC PROBLEM

- Relativistic field: collection of harmonic oscill's (scalar field)

$$E_0 = \frac{\hbar c}{2} \sum_n \omega_n, \quad \omega = k^2 + m^2/\hbar^2, \quad k = 2\pi/\lambda$$

- Evaluating in a box and putting a cut-off at maximum k_{max} corresp'ng to QFT physics (e.g., Planck energy)

$$M_P/M_{ew} \sim 10^{16}, \quad M_P/M_{cc} \sim 10^{31}, \quad \rho \sim \frac{\hbar k_{Planck}^4}{16\pi^2} \sim 10^{123} \rho_{obs}$$

a thick **aether!** **Caldwell, Carroll** but **Gómez, Dvali:** species $\downarrow 10^{30}$

- **Observational tests** see nothing (or **very little**) of it:

\implies **(new) cosmological constant problem**

CC PROBLEM

- Relativistic field: collection of harmonic oscill's (scalar field)

$$E_0 = \frac{\hbar c}{2} \sum_n \omega_n, \quad \omega = k^2 + m^2/\hbar^2, \quad k = 2\pi/\lambda$$

- Evaluating in a box and putting a cut-off at maximum k_{max} corresp'ng to QFT physics (e.g., Planck energy)

$$M_P/M_{ew} \sim 10^{16}, \quad M_P/M_{cc} \sim 10^{31}, \quad \rho \sim \frac{\hbar k_{Planck}^4}{16\pi^2} \sim 10^{123} \rho_{obs}$$

a thick **aether!** **Caldwell, Carroll** but **Gómez, Dvali:** species $\downarrow 10^{30}$

- **Observational tests** see nothing (or **very little**) of it:

\implies **(new) cosmological constant problem**

- Very difficult to solve and we **do not** address this question directly
[Baum, Hawking, Coleman, Polchinsky, Weinberg,...]

CC PROBLEM

- Relativistic field: collection of harmonic oscill's (scalar field)

$$E_0 = \frac{\hbar c}{2} \sum_n \omega_n, \quad \omega = k^2 + m^2/\hbar^2, \quad k = 2\pi/\lambda$$

- Evaluating in a box and putting a cut-off at maximum k_{max} corresp'ng to QFT physics (e.g., Planck energy)

$$M_P/M_{ew} \sim 10^{16}, \quad M_P/M_{cc} \sim 10^{31}, \quad \rho \sim \frac{\hbar k_{Planck}^4}{16\pi^2} \sim 10^{123} \rho_{obs}$$

a thick **aether!** **Caldwell, Carroll** but **Gómez, Dvali**: species $\downarrow 10^{30}$

- **Observational tests** see nothing (or **very little**) of it:

\implies **(new) cosmological constant problem**

- Very difficult to solve and we **do not** address this question directly

[Baum, Hawking, Coleman, Polchinsky, Weinberg,...]

- What we **do consider** —with relative success in some different approaches— is the **additional** contribution to the cc coming from the **non-trivial topology** of space or from specific **boundary conditions** imposed on braneworld models:

\implies **kind of cosmological Casimir effect**

Cosmo-Topol Casimir Eff't & Alternat's

- A. Assuming one is able to prove that the ground value of the cc is zero [Dolgov 1983; Ford 1987, 2002; Tsamis & Woodard 1998]
→ left with this incremental value coming from the topology or BCs

Cosmo-Topol Casimir Eff't & Alternat's

- A. Assuming one is able to prove that the ground value of the cc is zero [Dolgov 1983; Ford 1987, 2002; Tsamis & Woodard 1998]
→ left with this incremental value coming from the topology or BCs
- We have shown (with different examples) that this value acquires the correct order of magnitude —corresponding to the one coming from the observed acceleration in the expansion of our universe— in some reasonable models involving:

Cosmo-Topol Casimir Eff't & Alternat's

- A. Assuming one is able to prove that the ground value of the cc is zero [Dolgov 1983; Ford 1987, 2002; Tsamis & Woodard 1998]
→ left with this incremental value coming from the topology or BCs
- We have shown (with different examples) that this value acquires the correct order of magnitude —corresponding to the one coming from the observed acceleration in the expansion of our universe— in some reasonable models involving:
 - (a) small and large compactified scales JPA39(06)6299

Cosmo-Topol Casimir Eff't & Alternat's

- A. Assuming one is able to prove that the ground value of the cc is zero [Dolgov 1983; Ford 1987, 2002; Tsamis & Woodard 1998]
→ left with this **incremental value** coming from the topology or BCs
- We have shown (with different examples) that this value acquires the **correct order of magnitude** —corresponding to the one coming from the observed acceleration in the expansion of our universe— in some reasonable models involving:
 - (a) small and large compactified scales [JPA39\(06\)6299](#)
 - (b) dS & AdS worldbranes [hep-th/0209242](#)

Cosmo-Topol Casimir Eff't & Alternat's

- A. Assuming one is able to prove that the ground value of the cc is zero [Dolgov 1983; Ford 1987, 2002; Tsamis & Woodard 1998]
→ left with this incremental value coming from the topology or BCs
- We have shown (with different examples) that this value acquires the correct order of magnitude —corresponding to the one coming from the observed acceleration in the expansion of our universe— in some reasonable models involving:
 - (a) small and large compactified scales JPA39(06)6299
 - (b) dS & AdS worldbranes hep-th/0209242
 - (c) supergraviton theo's (discret dims, deconstr) hep-th/0312269

Cosmo-Topol Casimir Eff't & Alternat's

- A. Assuming one is able to prove that the ground value of the cc is zero [Dolgov 1983; Ford 1987, 2002; Tsamis & Woodard 1998]
→ left with this incremental value coming from the topology or BCs
- We have shown (with different examples) that this value acquires the correct order of magnitude —corresponding to the one coming from the observed acceleration in the expansion of our universe— in some reasonable models involving:
 - (a) small and large compactified scales JPA39(06)6299
 - (b) dS & AdS worldbranes hep-th/0209242
 - (c) supergraviton theo's (discret dims, deconstr) hep-th/0312269
- B. Other alternatives: (i) L Faddeev 0911.0282 (Adler '82)
Newton const in E-H Lag has dim of mass → non-renormalizability
Describe gravity by vector field (as Higgs mechanism)
(ii) Porto & Zee 0910.3716 Dynamical critical behavior of gravity in euIR sector and a mechanism to relax the cc. Also Shapiro+Sola, ...

More recent alternatives (a sample)

- (iii) E Mottola 1006.3567 Effective field theory approach
 - Casimir effect in flat s-t and large quantum backreaction are effects at the horizon scale of cosmological s-t
 - imply the cosmological VE is dynamical
 - its value depends on macroscopic BCs at the cosm horizon scale, rather than on the extreme ultraviolet Planck scale [we, on both BCs]

More recent alternatives (a sample)

- (iii) E Mottola 1006.3567 Effective field theory approach
 - Casimir effect in flat s-t and large quantum backreaction are effects at the horizon scale of cosmological s-t
 - imply the cosmological VE is dynamical
 - its value depends on macroscopic BCs at the cosm horizon scale, rather than on the extreme ultraviolet Planck scale [we, on both BCs]
- (iv) T Padmanabhan Ad Sci Lett 2 74 09 cc problem and explaining DE independent issues: first find mechanism to make the cc vanish
 - new degrees of freedom, kind of ‘gauge freedom’
 - to absorb any λ while maintaining general covariance
 - could succeed in making gravity decouple from the bulk VE
 - emergent gravity approach: thermodynamic description is far more general than just Einstein theory
 - observed cc should be a relic of quantum gravitational physics and arise from degrees of freedom which scale as the surface area
 - numerics: $L_\Lambda/L_P \sim \exp \sqrt{2} \pi^4 \sim 10^{60}$ (hierarchy squared) $\sim 10^{61}$

- (v) Shao & Chen 1005.1920 no attempt at explaining the old cc prob
 - an extremely small quantum correction can in fact be produced quite naturally from a massive bulk field, introducing a massive bulk fermion
 - naturally as superpartner of the radion field in a SUSY theory (especially the string theory realization) of brane-world scenario
 - in particle physics Grossman and Neubert used a massive bulk fermion to understand the neutrino mass hierarchy
 - use Goldberger-Wise mechanism where massive bulk scalar field with brane self-interaction induces stabilizing potential
 - could overwhelm the small fermionic Casimir energy & sign ?

- (v) [Shao & Chen 1005.1920](#) no attempt at explaining the old cc prob
 - an extremely small quantum correction can in fact be produced quite naturally from a massive bulk field, introducing a massive bulk fermion
 - naturally as superpartner of the radion field in a SUSY theory (especially the string theory realization) of brane-world scenario
 - in particle physics Grossman and Neubert used a massive bulk fermion to understand the neutrino mass hierarchy
 - use Goldberger-Wise mechanism where massive bulk scalar field with brane self-interaction induces stabilizing potential
 - could overwhelm the small fermionic Casimir energy & sign ?
- (vi) [JA Dixon 1006.2334](#) CyberSUSY solves the cc problem
 - a new mechanism for SUSY breaking
 - its realization mixes elementary and composite states
 - SUSY anomalies present, generates spectrum for SUSY breaking consistent with known particles
 - no cc generated, because SUSY is not spontaneously broken ...

The Braneworld Case

1. Braneworld may help to solve:

- the hierarchy problem
- the cosmological constant problem

2. Presumably, the bulk Casimir effect will play a role in the construction (radion stabilization) of braneworlds [A Flachi]

- Bulk Casimir effect (effective potential) for a conformal or massive scalar field
- Bulk is a 5-dim AdS or dS space with 2/1 4-dim dS brane (our universe)
- Consistent with observational data even for relatively large extra dimension

Previous work:

- flat space brane
- bulk conformal scalar field
- conclusion: no CE

We used **zeta regularization** at full power, with **positive** results!

EE, Nojiri, Odintsov, Ogushi, PRD67(2003)063515, hep-th/0209242 *Casimir effect in de Sitter and Anti-de Sitter braneworlds* EE, Odintsov, Saharian PRD79(2009)065023, 0902.0717 *Repulsive Casimir effect from extra dimensions and Robin BC: from branes to pistons*

The Sign of the Casimir Force

- Many papers dealing on this issue: here just short account

The Sign of the Casimir Force

- Many papers dealing on this issue: here just short account
- Casimir calculation: **attractive** force

The Sign of the Casimir Force

- Many papers dealing on this issue: here just short account
- Casimir calculation: **attractive** force
- Boyer got **repulsion** [TH, **Phys Rev**, 174 (1968)] for a spherical shell. It is a special case requiring stringent material properties of the sphere and a perfect geometry and BC

The Sign of the Casimir Force

- Many papers dealing on this issue: here just short account
- Casimir calculation: **attractive** force
- Boyer got **repulsion** [TH, *Phys Rev*, 174 (1968)] for a spherical shell. It is a special case requiring stringent material properties of the sphere and a perfect geometry and BC
- Systematic calculation, for different fields, BCs, and dimensions
J Ambjørn, S Wolfram, *Ann Phys NY* 147, 1 (1983) **attract, repuls**

The Sign of the Casimir Force

- Many papers dealing on this issue: here just short account
- Casimir calculation: [attractive](#) force
- Boyer got [repulsion](#) [TH, [Phys Rev, 174 \(1968\)](#)] for a spherical shell. It is a special case requiring stringent material properties of the sphere and a perfect geometry and BC
- Systematic calculation, for different fields, BCs, and dimensions
[J Ambjørn, S Wolfram, Ann Phys NY 147, 1 \(1983\)](#) [attract, repuls](#)
- Possibly not relevant at lab scales, but very important for cosmological models

The Sign of the Casimir Force

- Many papers dealing on this issue: here just short account
- Casimir calculation: **attractive** force
- Boyer got **repulsion** [TH, *Phys Rev*, 174 (1968)] for a spherical shell. It is a special case requiring stringent material properties of the sphere and a perfect geometry and BC
- Systematic calculation, for different fields, BCs, and dimensions
J Ambjørn, S Wolfram, Ann Phys NY 147, 1 (1983) **attract, repuls**
- Possibly not relevant at lab scales, but very important for cosmological models
- More general results: **Kenneth, Klich, PRL 97, 160401 (2006)**
a mirror pair of dielectric bodies always attract each other
CP Bachas, J Phys A40, 9089 (2007) from a general property of
Euclidean QFT '**reflection positivity**' (Osterwalder - Schrader 73, 75):
∃ of positive Hilbert space and self-adjoint non-negative Hamiltonian

- E.g. \exists correlation inequality: $\langle f\Theta(f) \rangle > 0$
 Θ reflection with respect to a 3-dim hyperplane in R^4
the action of Θ on f is anti-unitary $\Theta(cf) = c^*\Theta(f)$

- E.g. \exists correlation inequality: $\langle f\Theta(f) \rangle > 0$
 Θ reflection with respect to a 3-dim hyperplane in R^4
the action of Θ on f is anti-unitary $\Theta(cf) = c^*\Theta(f)$
- The existence of the reflection operator Θ is a consequence of unitarity only, and makes no assumptions about the discrete C, P, T symmetries

- E.g. \exists correlation inequality: $\langle f\Theta(f) \rangle > 0$
 Θ reflection with respect to a 3-dim hyperplane in R^4
the action of Θ on f is anti-unitary $\Theta(cf) = c^*\Theta(f)$
- The existence of the reflection operator Θ is a consequence of unitarity only, and makes no assumptions about the discrete C, P, T symmetries
- Boyer's result does not contradict the theorem, since cutting an elastic shell into two rigid hemispheres is a mathematically singular operation (which introduces divergent edge contributions)

- E.g. \exists correlation inequality: $\langle f\Theta(f) \rangle > 0$
 Θ reflection with respect to a 3-dim hyperplane in R^4
the action of Θ on f is anti-unitary $\Theta(cf) = c^*\Theta(f)$
- The existence of the reflection operator Θ is a consequence of unitarity only, and makes no assumptions about the discrete C, P, T symmetries
- Boyer's result does not contradict the theorem, since cutting an elastic shell into two rigid hemispheres is a mathematically singular operation (which introduces divergent edge contributions)
- Theorem does not apply for

- E.g. \exists correlation inequality: $\langle f\Theta(f) \rangle > 0$
 Θ reflection with respect to a 3-dim hyperplane in R^4
the action of Θ on f is anti-unitary $\Theta(cf) = c^*\Theta(f)$
- The existence of the reflection operator Θ is a consequence of unitarity only, and makes no assumptions about the discrete C, P, T symmetries
- Boyer's result does not contradict the theorem, since cutting an elastic shell into two rigid hemispheres is a mathematically singular operation (which introduces divergent edge contributions)
- Theorem does not apply for
 - mirror probes in a Fermi sea (chemical-potential term), eg when electron-gas fluctuations become important

- E.g. \exists correlation inequality: $\langle f\Theta(f) \rangle > 0$
 Θ reflection with respect to a 3-dim hyperplane in R^4
the action of Θ on f is anti-unitary $\Theta(cf) = c^*\Theta(f)$
- The existence of the reflection operator Θ is a consequence of unitarity only, and makes no assumptions about the discrete C, P, T symmetries
- Boyer's result does not contradict the theorem, since cutting an elastic shell into two rigid hemispheres is a mathematically singular operation (which introduces divergent edge contributions)
- Theorem does not apply for
 - mirror probes in a Fermi sea (chemical-potential term), eg when electron-gas fluctuations become important
 - periodic BCs for fermions

- E.g. \exists correlation inequality: $\langle f\Theta(f) \rangle > 0$
 Θ reflection with respect to a 3-dim hyperplane in R^4
the action of Θ on f is anti-unitary $\Theta(cf) = c^*\Theta(f)$
- The existence of the reflection operator Θ is a consequence of unitarity only, and makes no assumptions about the discrete C, P, T symmetries
- Boyer's result does not contradict the theorem, since cutting an elastic shell into two rigid hemispheres is a mathematically singular operation (which introduces divergent edge contributions)
- Theorem does not apply for
 - mirror probes in a Fermi sea (chemical-potential term), eg when electron-gas fluctuations become important
 - periodic BCs for fermions
 - Robin BCs in general \Leftarrow

Casimir eff in brworld's w large extra dim

- Casimir energy for massive scalar field with an arbitrary curvature coupling, obeying Robin BCs on two codim-1 parallel plates embedded in background spacetime $R^{(D_1-1,1)} \times \Sigma$, Σ compact internal space

Casimir eff in brworld's w large extra dim

- Casimir energy for massive scalar field with an arbitrary curvature coupling, obeying Robin BCs on two codim-1 parallel plates embedded in background spacetime $R^{(D_1-1,1)} \times \Sigma$, Σ compact internal space
- Most general case: constants in the BCs different for the two plates
It is shown that Robin BCs with different coefficients are necessary to obtain repulsive Casimir forces

Casimir eff in brworld's w large extra dim

- Casimir energy for massive scalar field with an arbitrary curvature coupling, obeying Robin BCs on two codim-1 parallel plates embedded in background spacetime $R^{(D_1-1,1)} \times \Sigma$, Σ compact internal space
- Most general case: constants in the BCs different for the two plates
It is shown that Robin BCs with different coefficients are necessary to obtain repulsive Casimir forces
- Robin type BCs are an extension of Dirichlet and Neumann's
⇒ most suitable to describe physically realistic situations

Casimir eff in brworld's w large extra dim

- Casimir energy for massive scalar field with an arbitrary curvature coupling, obeying Robin BCs on two codim-1 parallel plates embedded in background spacetime $R^{(D_1-1,1)} \times \Sigma$, Σ compact internal space
- Most general case: constants in the BCs different for the two plates
It is shown that Robin BCs with different coefficients are necessary to obtain repulsive Casimir forces
- Robin type BCs are an extension of Dirichlet and Neumann's
⇒ most suitable to describe physically realistic situations
- Genuinely appear in: → vacuum effects for a confined charged scalar field in external fields [Ambjørn ea 83],
→ spinor and gauge field theories,
→ quantum gravity and supergravity [Luckock ea 91]
Can be made conformally invariant, purely-Neumann conditions cannot
⇒ needed for conformally invariant theories with BC, to preserve cf invar

Casimir eff in brworld's w large extra dim

- Casimir energy for massive scalar field with an arbitrary curvature coupling, obeying Robin BCs on two codim-1 parallel plates embedded in background spacetime $R^{(D_1-1,1)} \times \Sigma$, Σ compact internal space
- Most general case: constants in the BCs different for the two plates
It is shown that Robin BCs with different coefficients are necessary to obtain repulsive Casimir forces
- Robin type BCs are an extension of Dirichlet and Neumann's
⇒ most suitable to describe physically realistic situations
- Genuinely appear in: → vacuum effects for a confined charged scalar field in external fields [Ambjørn ea 83],
→ spinor and gauge field theories,
→ quantum gravity and supergravity [Luckock ea 91]
Can be made conformally invariant, purely-Neumann conditions cannot
⇒ needed for conformally invariant theories with BC, to preserve cf invar
- Quantum scalar field with Robin BCs on boundary of cavity violates Bekenstein's entropy-to-energy bound near certain points in the space of the parameter defining the boundary condition [Solodukhin 01]

Gravity Eqs as Eqs of State: $f(R)$ Case

- The cosmological constant as an “integration constant”

T. Padmanabhan; D. Blas, J. Garriga, E. Alvarez ...

Unimodular Gravity

Also I Shapiro, J Solà,... cc RG flow

Gravity Eqs as Eqs of State: $f(R)$ Case

- The cosmological constant as an “integration constant”
T. Padmanabhan; D. Blas, J. Garriga, E. Alvarez ...
Unimodular Gravity Also I Shapiro, J Solà,... cc RG flow
- Ted Jacobson [PRL1995] obtained Einstein's equations from
local thermodynamics arguments only

Gravity Eqs as Eqs of State: $f(R)$ Case

- The cosmological constant as an “integration constant”
T. Padmanabhan; D. Blas, J. Garriga, E. Alvarez ...
Unimodular Gravity Also I Shapiro, J Solà,... cc RG flow
- Ted Jacobson [PRL1995] obtained Einstein’s equations from
local thermodynamics arguments only
- By way of generalizing black hole thermodynamics to
space-time thermodynamics as seen by a local observer

Gravity Eqs as Eqs of State: $f(R)$ Case

- The cosmological constant as an “integration constant”
T. Padmanabhan; D. Blas, J. Garriga, E. Alvarez ...
Unimodular Gravity Also I Shapiro, J Solà,... cc RG flow
- Ted Jacobson [PRL1995] obtained Einstein’s equations from
local thermodynamics arguments only
- By way of generalizing black hole thermodynamics to
space-time thermodynamics as seen by a local observer
- This strongly suggests, in a fundamental context:
Einstein’s Eqs are to be viewed as EoS

Gravity Eqs as Eqs of State: $f(R)$ Case

- The cosmological constant as an “integration constant”
T. Padmanabhan; D. Blas, J. Garriga, E. Alvarez ...
Unimodular Gravity Also I Shapiro, J Solà,... cc RG flow
- Ted Jacobson [PRL1995] obtained Einstein’s equations from local thermodynamics arguments only
- By way of generalizing black hole thermodynamics to space-time thermodynamics as seen by a local observer
- This strongly suggests, in a fundamental context:
Einstein’s Eqs are to be viewed as EoS
- Should, probably, not be taken as basic for quantizing gravity

Gravity Eqs as Eqs of State: $f(R)$ Case

- The cosmological constant as an “integration constant”
T. Padmanabhan; D. Blas, J. Garriga, E. Alvarez ...
Unimodular Gravity Also I Shapiro, J Solà,... cc RG flow
- Ted Jacobson [PRL1995] obtained Einstein’s equations from local thermodynamics arguments only
- By way of generalizing black hole thermodynamics to space-time thermodynamics as seen by a local observer
- This strongly suggests, in a fundamental context:
Einstein’s Eqs are to be viewed as EoS
- Should, probably, not be taken as basic for quantizing gravity
- C. Eling, R. Guedens, T. Jacobson [PRL2006]: extension to polynomial $f(R)$ gravity but as non-equilibrium thermodyn.
Also Erik Verlinde (private discussions)

● **Jacobson's argument:** basic thermodynamic relation

$$\delta Q = T\delta S$$

- entropy proportional to variation of the horizon area: $\delta S = \eta \delta \mathcal{A}$
- local temperature T defined as **Unruh temp**: $T = \hbar k / 2\pi$
- functional dependence of S wrt energy and size of system

- **Jacobson's argument:** basic thermodynamic relation

$$\delta Q = T \delta S$$

- entropy proportional to variation of the horizon area: $\delta S = \eta \delta \mathcal{A}$
 - local temperature T defined as **Unruh temp**: $T = \hbar k / 2\pi$
 - functional dependence of S wrt energy and size of system
- **Key point in our generalization:** the definition of the local entropy (Iyer+Wald 93: local boost inv, Noether charge)

$$S = -2\pi \int_{\Sigma} E_R^{pqrs} \epsilon_{pq} \epsilon_{rs}, \quad \delta S = \delta (\eta_e A)$$

η_e is a function of the metric and its derivatives to a given order

$$\eta_e = \eta_e \left(g_{ab}, R_{cdef}, \nabla^{(l)} R_{pqrs} \right)$$

- **Jacobson's argument:** basic thermodynamic relation

$$\delta Q = T \delta S$$

- entropy proportional to variation of the horizon area: $\delta S = \eta \delta \mathcal{A}$
 - local temperature T defined as **Unruh temp**: $T = \hbar k / 2\pi$
 - functional dependence of S wrt energy and size of system
- **Key point in our generalization:** the definition of the local entropy (Iyer+Wald 93: local boost inv, Noether charge)

$$S = -2\pi \int_{\Sigma} E_R^{pqrs} \epsilon_{pq} \epsilon_{rs}, \quad \delta S = \delta (\eta_e A)$$

η_e is a function of the metric and its derivatives to a given order

$$\eta_e = \eta_e \left(g_{ab}, R_{cdef}, \nabla^{(l)} R_{pqrs} \right)$$

- **Case of $\mathbf{f}(R)$ gravities:** $\mathbf{L} = \mathbf{f}(R, \nabla^n R)$

- Also the concept of an **effective Newton constant** for graviton exchange (**effective propagator**)

$$\begin{aligned}\frac{1}{8\pi G_{eff}} &= E_R^{pqrs} \epsilon_{pq} \epsilon_{rs} = \frac{\partial \mathbf{f}}{\partial R} (g^{pr} g^{qs} - g^{qr} g^{ps}) \epsilon_{pq} \epsilon_{rs} \\ &= \frac{\partial \mathbf{f}}{\partial R} = \frac{\eta_e}{2\pi}, \quad S = \frac{A}{4 G_{eff}}\end{aligned}$$

- Also the concept of an **effective Newton constant** for graviton exchange (**effective propagator**)

$$\begin{aligned} \frac{1}{8\pi G_{eff}} &= E_R^{pqrs} \epsilon_{pq} \epsilon_{rs} = \frac{\partial \mathbf{f}}{\partial R} (g^{pr} g^{qs} - g^{qr} g^{ps}) \epsilon_{pq} \epsilon_{rs} \\ &= \frac{\partial \mathbf{f}}{\partial R} = \frac{\eta_e}{2\pi}, \quad S = \frac{A}{4 G_{eff}} \end{aligned}$$

- For these theories, the different polarizations of the gravitons only enter in the definition of the **effective Newton constant through the metric itself**

- Also the concept of an **effective Newton constant** for graviton exchange (**effective propagator**)

$$\begin{aligned} \frac{1}{8\pi G_{eff}} &= E_R^{pqrs} \epsilon_{pq} \epsilon_{rs} = \frac{\partial \mathbf{f}}{\partial R} (g^{pr} g^{qs} - g^{qr} g^{ps}) \epsilon_{pq} \epsilon_{rs} \\ &= \frac{\partial \mathbf{f}}{\partial R} = \frac{\eta_e}{2\pi}, \quad S = \frac{A}{4 G_{eff}} \end{aligned}$$

- For these theories, the different polarizations of the gravitons only enter in the definition of the **effective Newton constant through the metric itself**
- Final result, for $\mathbf{f}(R)$ gravities:
the local field equations can be thought of as an equation of state of equilibrium thermodynamics (as in the GR case)

- Jacobson's argum **non-trivially extended to $f(R)$** gravity field eqs as EoS of local space-time thermodynamics
EE, P. Silva, Phys Rev D78, 061501(R) (2008), arXiv:0804.3721v2

- Jacobson's argument **non-trivially extended to $f(R)$** gravity field eqs as EoS of local space-time thermodynamics
EE, P. Silva, Phys Rev D78, 061501(R) (2008), arXiv:0804.3721v2
- By means of a **more general** definition of local entropy, using **Wald's definition** of dynamic BH entropy
RM Wald PRD1993; V Iyer, RM Wald PRD1994

- Jacobson's argum **non-trivially extended to $f(R)$** gravity field eqs as EoS of local space-time thermodynamics
EE, P. Silva, Phys Rev D78, 061501(R) (2008), arXiv:0804.3721v2
- By means of a **more general** definition of local entropy, using **Wald's definition** of dynamic BH entropy
RM Wald PRD1993; V Iyer, RM Wald PRD1994
- And also the concept of an **effective Newton constant** for graviton exchange (effective propagator)
R. Brustein, D. Gorbonos, M. Hadad, arXiv:0712.3206

- Jacobson's argum **non-trivially extended to $f(R)$** gravity field eqs as EoS of local space-time thermodynamics
EE, P. Silva, Phys Rev D78, 061501(R) (2008), arXiv:0804.3721v2
- By means of a **more general** definition of local entropy, using **Wald's definition** of dynamic BH entropy
RM Wald PRD1993; V Iyer, RM Wald PRD1994
- And also the concept of an **effective Newton constant** for graviton exchange (effective propagator)
R. Brustein, D. Gorbonos, M. Hadad, arXiv:0712.3206
- S-F Wu, G-H Yang, P-M Zhang, arXiv:0805.4044, **direct extension** of our results to **Brans-Dicke** and **scalar-tensor** gravities
T Zhu, Ji-R Ren and S-F Mo, arXiv:0805.1162 [gr-qc];
C Eling, arXiv:0806.3165 [hep-th]; R-G Cai, L-M Cao and Y-P Hu, arXiv:0807.1232 [hep-th] & arXiv:0809.1554 [hep-th]

- Jacobson's argument **non-trivially extended to $f(R)$** gravity field eqs as EoS of local space-time thermodynamics
EE, P. Silva, Phys Rev D78, 061501(R) (2008), arXiv:0804.3721v2
- By means of a **more general** definition of local entropy, using **Wald's definition** of dynamic BH entropy
RM Wald PRD1993; V Iyer, RM Wald PRD1994
- And also the concept of an **effective Newton constant** for graviton exchange (effective propagator)
R. Brustein, D. Gorbonos, M. Hadad, arXiv:0712.3206
- S-F Wu, G-H Yang, P-M Zhang, arXiv:0805.4044, **direct extension** of our results to **Brans-Dicke** and **scalar-tensor** gravities
T Zhu, Ji-R Ren and S-F Mo, arXiv:0805.1162 [gr-qc];
C Eling, arXiv:0806.3165 [hep-th]; R-G Cai, L-M Cao and Y-P Hu, arXiv:0807.1232 [hep-th] & arXiv:0809.1554 [hep-th]

DANKE FÜR IHRE AUFMERKSAMKEIT!